

Representation theory ^{of group th.} [relevance to physics]

We may represent group elements explicitly by matrices

$$\left. \begin{array}{l} g_1 \rightarrow \mathcal{D}(g_1) \quad \text{an } N \times N \text{ matrix} \\ g_2 \rightarrow \mathcal{D}(g_2) \quad \text{another} \\ \dots \\ g_n \rightarrow \mathcal{D}(g_n) \quad \text{an } n^{\text{th}} \text{ such matrix.} \end{array} \right\} \begin{array}{l} \mathcal{D}(g_1 g_2) = \mathcal{D}(g_1) \mathcal{D}(g_2) \\ \mathcal{D}(g_i^{-1}) = [\mathcal{D}(g_i)]^{-1} \\ \mathcal{D}(e) = I \end{array}$$

This set of matrices $\{\mathcal{D}(g_1), \dots, \mathcal{D}(g_n)\}$ is called a representation of the group G . They have the same multiplication ^{table} as the group elements. _{under matrix mult.}

One rep $\rightarrow \mathcal{D}(g_i) = 1 \forall i$. "Trivial representation." 1 (1-dimensional)

$$g_{m_1} \times g_{m_2} = g_{m_3} \rightarrow 1 \times 1 = 1 \checkmark$$

A faithful representation is one in which each group element is mapped into a different matrix, so we can infer the group multiplication table from the set $\{\mathcal{D}(g)\}$. The trivial rep. is clearly not faithful.

Faithful rep = isomorphism between group G and matrix group $\{\mathcal{D}(g)\}$
 Not faithful, = homomorphism " " " " " " " "

Representations typically arise in physics by looking at the effect of a transformation on the configuration space. Specific example of parity transformation in 1D operating on wavefunctions;

$$e\psi(x) = \psi(x) \quad P\psi(x) = \psi(-x) \quad P^2 = e$$

$$e^2 = e \quad eP = P \quad Pe = P$$

If we write the wavefunction as a column vector,

$$\begin{bmatrix} \psi(x) \\ \psi(-x) \end{bmatrix},$$

then we have

$$e \begin{bmatrix} \psi(x) \\ \psi(-x) \end{bmatrix} = \begin{bmatrix} \psi(x) \\ \psi(-x) \end{bmatrix} = \underbrace{\begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}}_{\mathcal{D}(e) = I} \begin{bmatrix} \psi(x) \\ \psi(-x) \end{bmatrix}$$

$$P \begin{bmatrix} \psi(x) \\ \psi(-x) \end{bmatrix} = \begin{bmatrix} \psi(-x) \\ \psi(x) \end{bmatrix} = \underbrace{\begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}}_{\mathcal{D}(P)} \begin{bmatrix} \psi(x) \\ \psi(-x) \end{bmatrix}$$

$$e^2 = e \quad \mathcal{D}(e)\mathcal{D}(e) = \begin{bmatrix} 1 & \\ & 1 \end{bmatrix} \begin{bmatrix} 1 & \\ & 1 \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} = \mathcal{D}(e) \checkmark$$

$$P^2 = e \quad \mathcal{D}(P)\mathcal{D}(P) = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} = \mathcal{D}(e) \checkmark$$

For a more general coordinate transformation $x' = Tx$ (e.g. 3D rotation, T_{ij} an orthog matrix) we may have a set of functions that mix among themselves under the transformation. e.g. $\psi_1(\vec{x}) = x$, $\psi_2(\vec{x}) = y$, $\psi_3(\vec{x}) = z$ only mix among themselves under a rotation.

Generally we write (for n such functions)

$$\underbrace{\mathcal{O}_T}_{\text{operation of } T \text{ on function } \psi_\mu} \psi_\mu(x) = \sum_{\nu=1}^n \mathcal{D}_{\nu\mu}^{(n)}(T) \psi_\nu(x) = \psi_\mu(T^{-1}x)$$

$$\mathcal{O}_T \psi_\mu(Tx) = \psi_\mu(x)$$

With this funny order of indices we have for the result of two successive operations

$$\begin{aligned}
 \mathcal{O}_{SR} \psi_\mu &= \mathcal{O}_S \mathcal{O}_R \psi_\mu = \mathcal{O}_S \sum_\nu \mathcal{D}_{\nu\mu}^{(n)}(R) \psi_\nu \\
 &= \sum_{\nu\sigma} \mathcal{D}_{\nu\mu}^{(n)}(R) \mathcal{D}_{\sigma\nu}^{(n)}(S) \psi_\sigma \\
 &= \sum_{\nu\sigma} [\mathcal{D}_{\sigma\nu}^{(n)}(S) \mathcal{D}_{\nu\mu}^{(n)}(R)] \psi_\sigma \\
 &= \sum_{\sigma} [\mathcal{D}^{(n)}(S) \cdot \mathcal{D}^{(n)}(R)]_{\sigma\mu} \psi_\sigma \\
 &= \sum_{\sigma} \mathcal{D}_{\sigma\mu}^{(n)}(SR) \psi_\sigma
 \end{aligned}$$

So the rep of a combined op $\mathcal{D}^{(n)}(SR) = \mathcal{D}^{(n)}(S) \cdot \mathcal{D}^{(n)}(R)$, just the conventional matrix product

In addition to the trivial rep $\mathcal{D}(g) = 1 \ \forall g$, two other reps are easy to write down. One is the "antisymmetric" rep,

$$\mathcal{D}(g) = \begin{cases} +1 & \text{if the order of } g \text{ is odd (smallest } m \ni g^m = e), \\ -1 & \text{" " " " " " " " even.} \end{cases}$$

To focus on specific e.g., recall S_3 with elements $e \ (12) \ (13) \ (23) \ (123) \ (132)$

trivial rep $\mathcal{D}(g) = 1 \ \forall g$

antisymmetric rep

$$\mathcal{D}(g) = \begin{cases} 1 & e, (123), (132) \text{ odd order (this det } e \text{ of order 1,} \\ -1 & (12), (13), (23) \text{ even order} \end{cases} \quad \begin{matrix} \text{Min } e^m = e \\ m \rightarrow m=1 \end{matrix}$$

The remaining "easy" rep. to write down is the "regular representation", which can be written directly from the group multiplication table.

The definition is

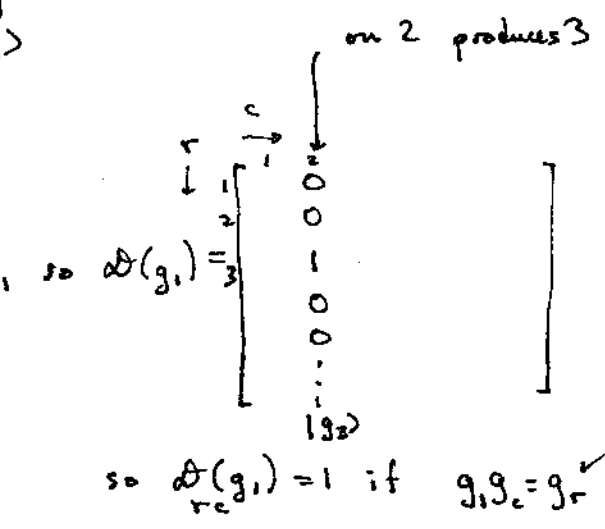
$$D_{rc}^{(n)}(g) = \begin{cases} 1 & \text{if } g g_o = g_r \\ 0 & \text{otherwise} \end{cases}$$

The way it works is to rep each group element as a column vector,

$$g_1 \rightarrow \begin{bmatrix} 1 \\ 0 \\ 0 \\ \vdots \\ 0 \end{bmatrix} \quad |g_1\rangle \quad g_2 \rightarrow \begin{bmatrix} 0 \\ 1 \\ 0 \\ \vdots \\ 0 \end{bmatrix} \quad |g_2\rangle \quad g_3 \rightarrow \begin{bmatrix} 0 \\ 0 \\ 1 \\ \vdots \\ 0 \end{bmatrix} \quad |g_3\rangle$$

so e.g. if $g_1 g_2 = g_3$,

$$D(g_1) \cdot \begin{bmatrix} 0 \\ 1 \\ 0 \\ \vdots \\ 0 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 1 \\ \vdots \\ 0 \end{bmatrix}$$



so $D_{rc}^{(n)}(g) = 1$ if $g_1 g_c = g_r$

Note the dimensionality of the reg rep = order of the group.

Note we can trivially change one rep to another by a similarity transformation. For any group elements $a, b, c \in \mathcal{G}$,

$$ab = c \xrightarrow{\text{rep}} D(a)D(b) = D(c)$$

Under a unitary transformation, we get a new set of matrices $\tilde{D}(g) = U^\dagger D(g) U$ $U^\dagger = U^{-1}$

note

$$\tilde{D}(a)\tilde{D}(b) = U^\dagger D(a) \underbrace{U U^\dagger}_I D(b) U = U^\dagger D(a)D(b)U = U^\dagger D(c)U = \tilde{D}(c)$$

So the new set $\{\tilde{D}(g)\}$ are also a rep.

Since a similarity transform is just a change of basis, ^{or coord sys} reps related by one are essentially identical and are called "equivalent".

Hence the matrices alone are not a very useful way of describing a representation. More useful way is to quote the "characters", which are the traces of the matrices,

$$\chi(g) \equiv \text{Tr} [D(g)].$$

This is invariant under a similarity transformation,

$$\chi(g) = \text{Tr} [D(g)],$$

$$\tilde{\chi}(g) = \text{Tr} [\tilde{D}(g)] = \text{Tr} [Y^\dagger D(g) Y] = \text{Tr} [D(g) \underbrace{Y Y^\dagger}_I] = \chi(g).$$

Note the $\chi(g)$ are class functions, if $g_1 = g g_2 g^{-1}$, then $\chi(g_1) = \chi(g_2)$.

A "character table" has one entry for each class.

Reducibility and Irreducible representations.

Given a set of matrices $\{D(g)\}$ that represent a group G , it may be possible to block-diagonalize them by a similarity transformation,

$$Y^\dagger D(g) Y = \begin{bmatrix} D_1(g) & & & \\ & D_2(g) & & \\ & & D_3(g) & \\ & & & D_4(g) \end{bmatrix} \quad \forall g, \text{ i.e.}$$

the rep $D(g)$ is really the sum of smaller matrices

which cannot be reduced any further in size.

We write this as

$$D(g) = D_1(g) \oplus D_2(g) \oplus D_3(g) \oplus D_4(g) \quad \text{and say}$$

$\rho(g)$ is the "direct sum" of $\rho_1, \rho_2, \rho_3, \rho_4$. (Any e.g. ρ_1 could appear once or not at all.)

A rep that can be block-diagonalized is referred to as "reducible", and the reps like ρ_1, \dots, ρ_4 that cannot be reduced to smaller reps are known as "irreducible" reps.

A general rep $\rho(g)$ can be decomposed into irreducible representations, and for a finite group the number of irreps is also finite.

Fund results (not proven) : 1) # of irreducible reps of $G = \#$ of classes κ

2) the dimensionalities $\{d_i\}$ of the matrices in the i^{th} irred rep satisfy

$$\sum_{i=1}^{\kappa} d_i^2 = n,$$

where n is the order of the group.

Specific example of S_3 : we learn from group mult table and looking for equivalent elements that S_3 has $\kappa=3$ classes, therefore it has 3 irreducible reps.

Since it has $S_n \rightarrow n!=6$ elements,

$$\sum_{i=1}^{\kappa=3} d_i^2 = n=6.$$

The only 3 ^{pos.} integers $\Rightarrow d_1^2 + d_2^2 + d_3^2 = 6$ are 1, 1, 2, so S_3 has two 1d irred reps (trivial and antisymmetric reps), and an irred rep of 2D matrices. (These 2D matrices are given explicitly by MW on p.433, table 16-1.)

Orthogonality relations of representations & characters.

Also derived from Schur's lemmas, see e.g. Hamermesh Ch 3.

Suppose we have two irreducible reps of a finite group G in terms of unitary matrices ($D^+ = D^{-1}$). (Can easily convert a nonunitary rep into a unitary one by a similarity transform; Wigner gives the transform explicitly Ch 9 Theorem I).

for irreducible representations

If the order of the group is n , the representation matrices satisfy

$$\sum_g D_{L_1 J_1}^{(r_1)}(g) D_{L_2 J_2}^{(r_2)}(g)^* = \frac{n}{d_1} \delta_{r_1 r_2} \delta_{L_1 L_2} \delta_{J_1 J_2}$$

dimension of r_1 matrices
must be same rep

Orthog theorem for group representation matrices.

Note this implies that the total # of irred reps of a finite group is itself finite: Can think of the $\{D_{ij}^{(r)}\}$ as a set of d_1^2 vectors

vec $\{ \underbrace{D_{11}^{(r_1)}(e)}_{\text{a number}}, \underbrace{D_{11}^{(r_1)}(g_1)}_{\text{a number}}, D_{11}^{(r_1)}(g_2), \dots, D_{11}^{(r_1)}(g_{n-1}) \} = \text{a } 1^{\text{st}} \text{ } n\text{-component vector}$

vec $\{ D_{d_1 d_1}^{(r_1)}(e), D_{d_1 d_1}^{(r_1)}(g_1), \dots, D_{d_1 d_1}^{(r_1)}(g_{n-1}) \} = \text{a } (d_1^2)^{\text{th}} \text{ } n\text{-component vector}$
 + sim for r_2

And all these vectors are orthogonal in this n -dimensional space.

Since you can have at most n orthog vectors in an n -dim space, we have

$$\sum_r d_r^2 \leq n$$

So even if all irreps were 1-dim you could have at most n of them.
 (This is the case for Abelian groups: all units, irreps. of an Abelian group are 1D.)

In fact this is an equality,

(Proof straightforward in Homework)
 (Uses decom. of reg. rep into irreps)

$$\boxed{\sum_r d_r^2 = n}$$

\nearrow Σ over all irreps
 \nearrow their dim ²
 \nearrow order of the group

This can sometimes be solved uniquely to specify the dims of irreps of a group.

We are really more interested in the characters of an irrep, since they are invariant under similarity transformation.

Taking the trace $\delta_{L_1, J_1} \delta_{L_2, J_2}$ of our representation matrix at the g relation, we have

$$\sum_g \chi^{(r_1)}(g) \chi^{(r_2)}(g)^* = \frac{n}{d_1} \delta_{r_1, r_2} \underbrace{\delta_{L_1, L_2} \delta_{L_1', L_2'}}_{\delta_{L_1, L_2} = d_1} = n \delta_{r_1, r_2}$$

or, since the characters $\chi(g)$ are the same for each class,

$$\boxed{\sum_k P_k \chi^{(r_1)}(C_k) \chi^{(r_2)}(C_k)^* = n \delta_{r_1, r_2}}$$

We can sometimes use these results to set up the "character table" for a group, which lists the characters $\chi^{(r_i)}(C_k)$ of each class C_k in each irrep r_i . Since the # classes = # irreps, it is a square table.

For the example of S_3 , we already know we have the identity or trivial rep

$$\rho(g) = 1 \quad \forall g \quad \rightarrow \chi(C_k) = 1 \quad \forall C_k$$

⊕ antisymm rep $\rho(g) = 1 \quad g = e, (123) \rightarrow \chi(C_k) = 1 \quad k=1,2$
 $\rho(g) = -1 \quad g = (12) \rightarrow \chi(C_k) = -1 \quad k=3$

e.g. ρ_k	# elems in class	class	rep		
			$\chi^{(1)}$	$\chi^{(2)}$	$\chi^{(3)}$
e	1	C_1	1	1	2
(123)	2	C_2	1	1	$\alpha = -1$
(12)	3	C_3	1	-1	$\beta = 0$

⊕ we know there is a 2D rep in addition, since

$$\sum_{r=1}^3 d_r^2 = n = 6$$

$$1^2 + 1^2 + d_3^2 = 6$$

$$\rightarrow d_3 = 2$$

For the 2 rep, $\rho(e) = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$, so $\chi(e) = 2$

The remaining elements in the character table can be found from the orthogonality relation for characters of irreps,

$$\sum_{k=1}^3 P_k \chi^{(r_1)}(C_k) \chi^{(r_2)}(C_k)^* = n \delta_{r_1, r_2}$$

$$\left. \begin{aligned} \chi_{r_1}^{(1)} = (1) \rightarrow 1 \cdot 1 \cdot 2 + 2 \cdot 1 \cdot \alpha^* + 3 \cdot 1 \cdot \beta^* &= 0 \\ \chi_{r_2}^{(3)} = (3) \end{aligned} \right\} \begin{aligned} \Sigma \neq \text{diff}, \alpha^* &= -1 \\ \beta^* &= 0 \end{aligned}$$

$$\left. \begin{aligned} \chi_{r_1}^{(2)} = (2) \rightarrow 1 \cdot 1 \cdot 2 + 2 \cdot 1 \cdot \alpha^* + 3 \cdot (-1) \cdot \beta^* &= 0 \\ \chi_{r_2}^{(3)} = (3) \end{aligned} \right\} \begin{aligned} \beta^* &= 0 \end{aligned}$$

n.b. $r_1 = (1) \quad r_2 = (2) \rightarrow 1 \cdot 1 \cdot 1 + 2 \cdot 1 \cdot 1 + 3 \cdot 1 \cdot (-1) = 0$ ✓ n.b. $r_1 = r_2 = (3) \quad 1 \cdot 2^2 + 2 \cdot (-1)^2 + 3 \cdot 0^2 = 6$ ✓ *nontrivial check*

Actually can think of $\frac{1}{\sqrt{n}} \left\{ \sqrt{p_k} x^{(r)}(c_k) \right\}_{k=1,2,\dots, N_{\text{classes}}=K}$

as one of a set of K orthonormal vectors in a K -dimensional space.

e.g. $1^{\text{st}} = \frac{1}{\sqrt{6}} (1, \sqrt{2}, \sqrt{3})$

$2^{\text{nd}} = \frac{1}{\sqrt{6}} (1, \sqrt{2}, -\sqrt{3})$

$3^{\text{rd}} = \frac{1}{\sqrt{6}} (2, -\sqrt{2}, 0)$

* you can check that these are an orthonormal set.