

Phy 512 review of topics (pre-midterm)

features of classical physics

atomic stability $W = \frac{2}{3} \frac{e^2 a^2}{c^3}$

discrete spectral lines

H atom $\rightarrow E_n \propto \frac{1}{n^2}$

UV catastrophe in black body rad.

Black body prob. (Planck)

modes, consisting of d.o.f.s $\frac{dN}{dk}, \frac{dN}{d\omega}$
in 1D 2D 3D

$E = kT$ each mode, $\infty \neq$ modes!

$\leftarrow P(\lambda) \propto \frac{1}{\lambda^4}$
 $\frac{dN}{d\lambda} =$ Rayleigh-Jeans law

checking
units !

Planck's assumption $E = \hbar\omega$ (steps in)
(1900)

$P(\lambda, T) \propto \frac{1}{\lambda^5 (e^{\frac{hc}{\lambda kT}} - 1)}$ ✓ cpt.

how to get relevant
 \int s and sums by
diff. w.r.t. params.

Bohr atom: classical except $L_n = n\hbar$. ($V = -\frac{e^2}{r}$,
 $e^2 = \hbar c \alpha$)

$r_n = n \cdot a_0$

$L \frac{\hbar^2}{me^2} = \frac{\hbar c}{m c} =$ Bohr radius

$E_n = -mc^2 \alpha^2 \cdot \frac{1}{2n^2}$

de Broglie "pilot waves"

(1924)

$p = \hbar k$

$E = \hbar\omega$

Davisson & Germer (1927)

Schrödinger's wave eqn. (1925)

$$-\frac{\hbar^2}{2m} \nabla^2 \Psi + V\Psi = i\hbar \frac{\partial}{\partial t} \Psi \quad \Psi(\vec{x}, t) \quad \text{TISE}$$

from $\frac{\vec{p}^2}{2m} + V(\vec{x}) = E$

$$-\frac{\hbar^2}{2m} \nabla^2 \psi + V\psi = E\psi \quad \Psi(\vec{x}, t) = e^{-iEt/\hbar} \psi(\vec{x})$$

TISE

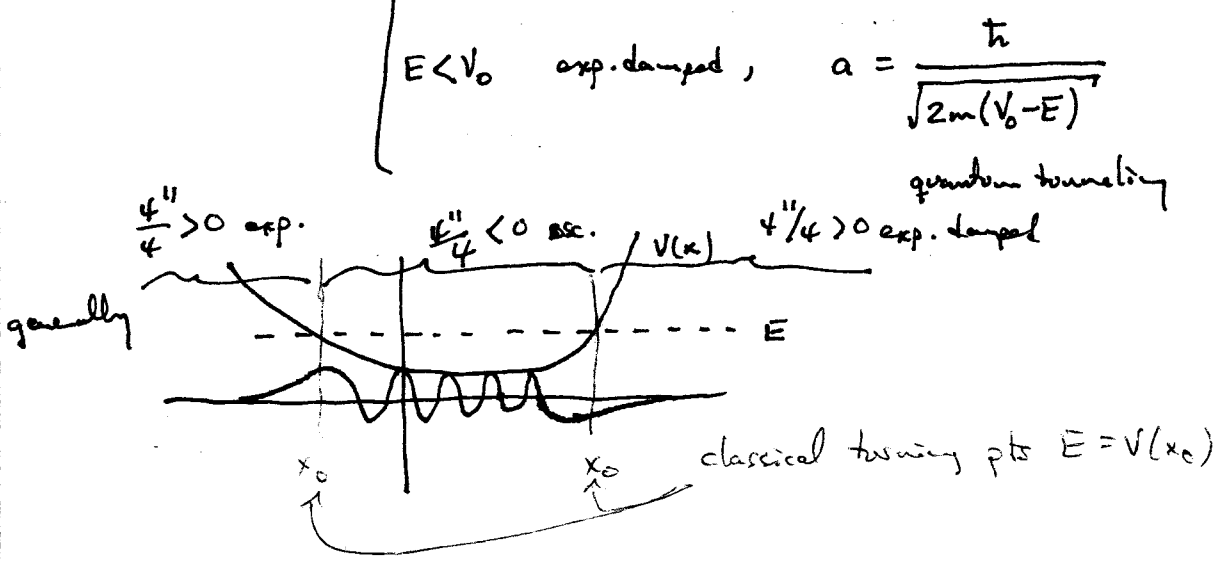
"energy eigenstates"

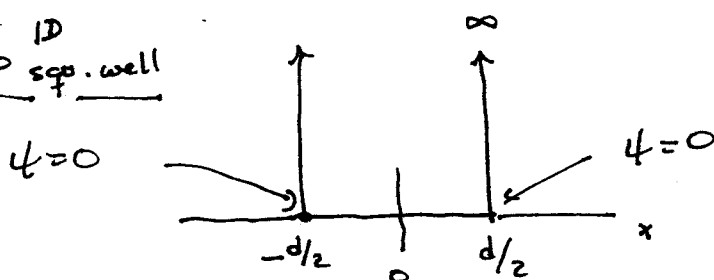
$|\Psi|^2$ & the "Copenhagen interpretation"

← energy eigenstate
1D problems

Const. V_0 $\left\{ \begin{array}{l} E > V_0: \quad \lambda = \frac{h}{p} = \frac{h}{\sqrt{2m(E-V_0)}}, \quad \pm k = \sqrt{2m(E-V_0)} \\ \\ E < V_0 \quad \text{exp. damped}, \quad a = \frac{\hbar}{\sqrt{2m(V_0-E)}} \end{array} \right.$

oscillatory region
classically allowed



Specific probs1D
 ∞ sq. well

$$E\psi = -\frac{\hbar^2}{2m}\psi''$$

even & odd (cos & sin) solns

$$\cos(k_n x), \sin(k_n x)$$

$$\lambda_n = \frac{2d}{n}$$

$$k_n = \frac{2\pi}{\lambda_n} = \frac{\pi n}{d}$$

$$\therefore E_n = n^2 \frac{\hbar^2 \pi^2}{2md^2}$$

$(N+1)^2$ if $N=0$ is the g.s.

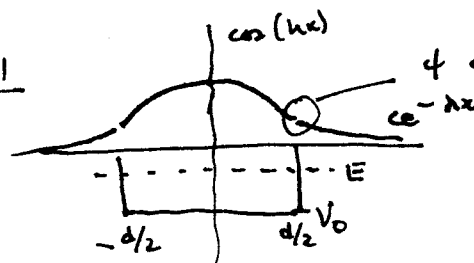
Normalization

$$\int |\psi|^2 dx = 1$$

Hers. unc. princ.

$\Delta x \cdot \Delta p_x \geq \hbar/2$ any wfn. (Gaussian gives min. uncertainty)

$$\Delta x = \left[\langle x^2 \rangle - \langle x \rangle^2 \right]^{1/2}$$

finite depth sq. well ψ & ψ' continuous.

to get E_n only,
matching ψ'/ψ
suffices.

(even solns) $x \tan x = \sqrt{R^2 - x^2}$

(odd) $\rightarrow -x \cot x$

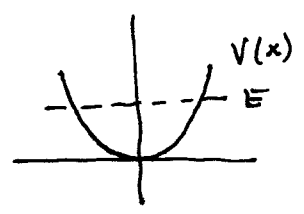
$$x = \frac{kd}{2}$$

$$R^2 = \frac{m|V_0|d^2}{2\hbar^2}$$

always have at least
1 b.s.

ID SHO

$$V(x) = \frac{1}{2} kx^2$$



$$-\frac{\hbar^2}{2m} \psi'' + \frac{1}{2} kx^2 \psi = E_n \psi$$

try $\psi = e^{-f}$

$$-\frac{\hbar^2}{2m} (-f'' + f'^2) + \frac{1}{2} kx^2 = E_0$$

↑ dom. large x

$$\psi_0 = \eta_0 e^{-\frac{\sqrt{km}}{2\hbar} x^2} \quad \checkmark \text{ exact soln} = \text{Gaussian}$$

excited states

$$\psi_n = p(x) e^{-\frac{\sqrt{km}}{2\hbar} x^2} = p(x) e^{-\frac{1}{2} cx^2} \quad c \equiv \frac{\sqrt{km}}{\hbar}$$

$s = \sqrt{c} x$ dimensionless

$$\frac{d^2 p}{ds^2} - 2s \frac{dp}{ds} + \left(\frac{E}{\frac{1}{2} \hbar \omega} - 1 \right) p = 0$$

2n

Hermite DE

solns are Hermite polynomials

$$\psi_n = \eta_n H_n(s) e^{-\frac{1}{2} s^2}$$

↳ $(-1)^n e^{s^2} \frac{d^n}{ds^n} e^{-s^2}$

gen fun. $F(r, s) = \sum_n \frac{r^n}{n!} H_n(s) = e^{2rs - r^2}$

orthog. $\int_{-\infty}^{\infty} H_n(s) H_m(s) e^{-s^2} ds = 2^n n! \sqrt{\pi} \delta_{nm}$

Why polynomials? (hence $\frac{E}{\frac{1}{2}\hbar\omega} - 1 = 2n$). normalizable solns,
localized near $r=0$.

norm $\eta_0 = \left(\frac{c}{\pi}\right)^{1/4}$

$$\eta_n = \frac{1}{\sqrt{2^n n!}} \left(\frac{c}{\pi}\right)^{1/4}$$

expectation values

$$\langle x \rangle = \int_{-\infty}^{\infty} x P(x) dx = \int_{-\infty}^{\infty} \psi^* x \psi dx \quad \text{"sandwich formula"}$$

↳ specify wfn like $\langle 0|x|0 \rangle = \int_{-\infty}^{\infty} \psi_0^*(x) x \psi_0(x) dx$

$$\langle 1|x|0 \rangle = \int_{-\infty}^{\infty} \psi_1^*(x) x \psi_0(x) dx$$

can also calc. expected values of operators, e.g. $p_x = -i\hbar \frac{d}{dx}$
 $p_x^2 = -\hbar^2 \frac{d^2}{dx^2}$

for the SHO we derived

$$\langle n|x^2|n \rangle = \frac{\hbar}{\sqrt{km}} \left(n + \frac{1}{2}\right) \text{ explicitly} \quad \left(\text{checks } \langle \psi|V(x)|\psi \rangle = \frac{1}{2} E_n\right).$$

SHO \pm ops

$$a = \sqrt{\frac{\hbar}{2m\omega}} \left(+ \frac{d}{dx} + \frac{m\omega}{\hbar} x \right)$$

$$a^\dagger = \sqrt{\frac{\hbar}{2m\omega}} \left(- \frac{d}{dx} + \frac{m\omega}{\hbar} x \right)$$

$$H = \hbar\omega \left(a^\dagger a + \frac{1}{2} \right)$$