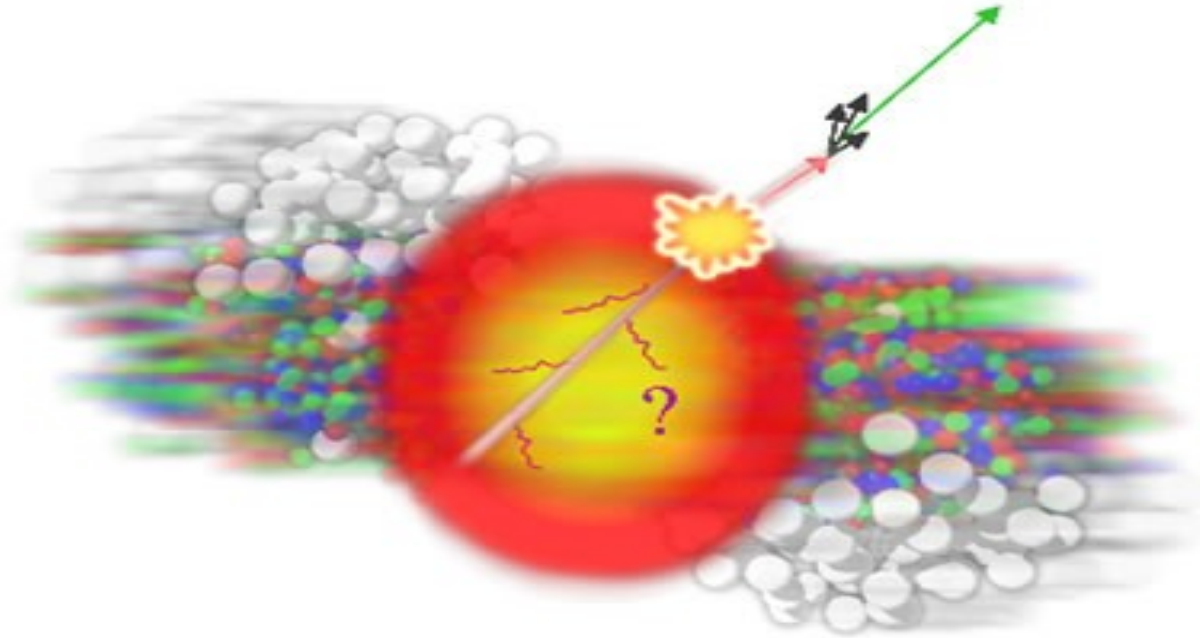


Jet spectra



Christine Nattrass

University of Tennessee, Knoxville

Questions an experimentalist should ask

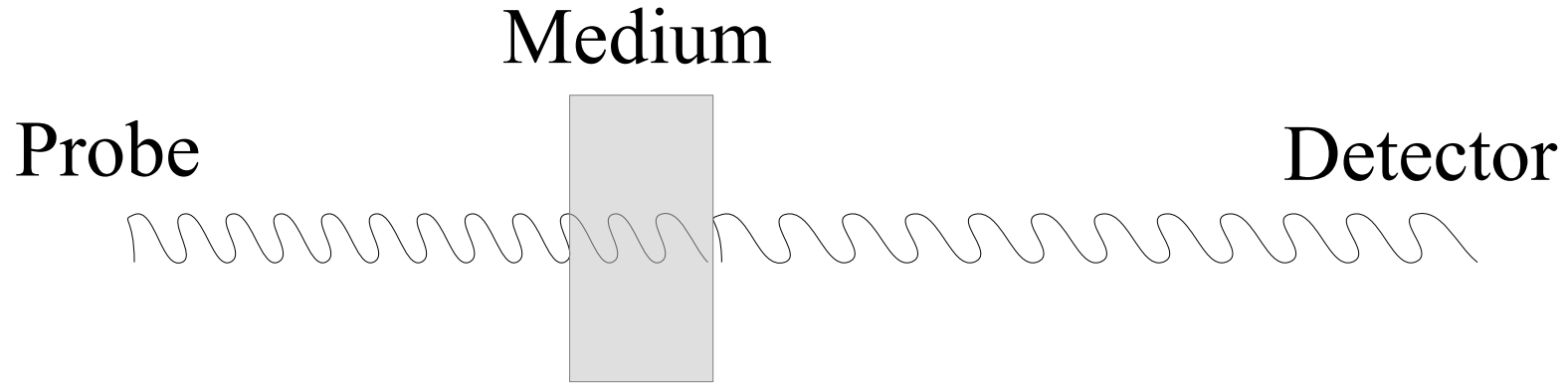
- What do I want to learn?
- What am I measuring?
- What assumptions am I making?
- What are the dominant uncertainties?
- How do I compare to models?

The answers for jets are highly non-trivial!

What do I want to learn? The cartoon picture



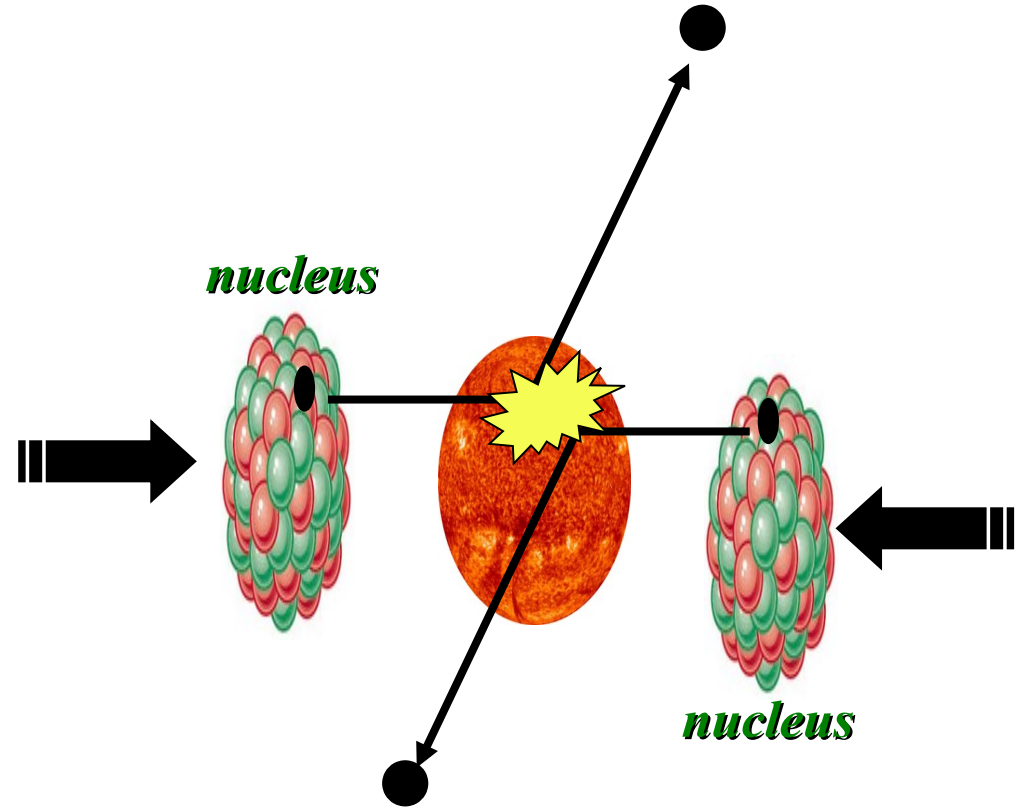
Probing the Quark Gluon Plasma



Want a probe which traveled through the collision
QGP is very short-lived ($\sim 1-10$ fm/c) \rightarrow
cannot use an external probe

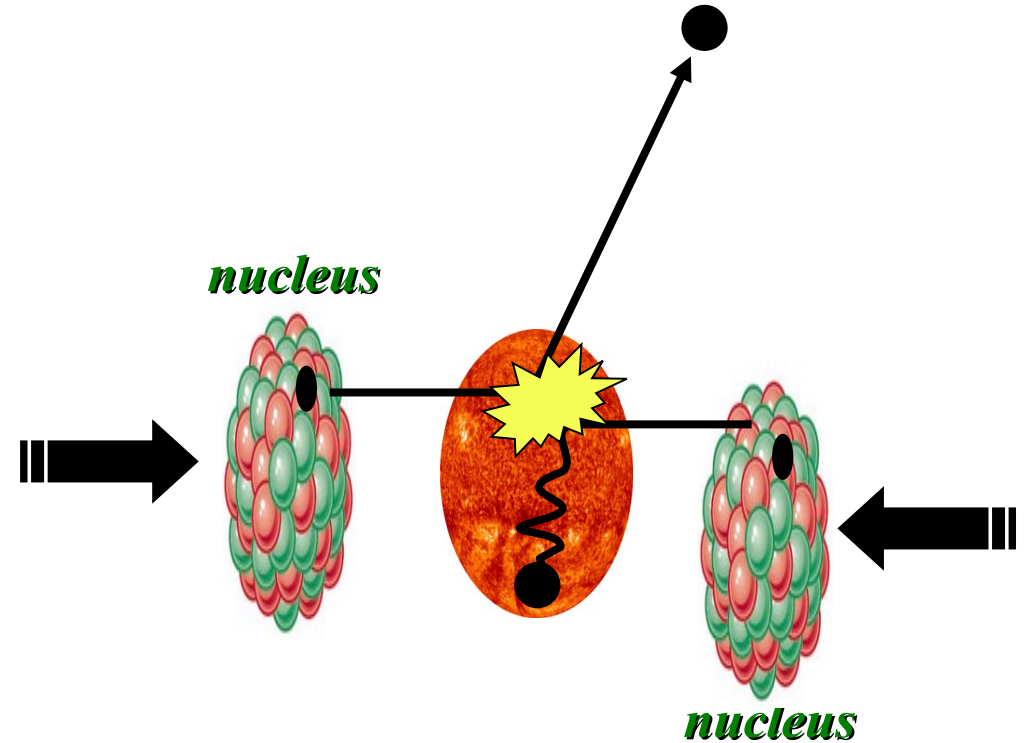
Probes of the Quark Gluon Plasma

Want a probe which traveled through the medium
QGP is short lived \rightarrow need a probe created in the collision



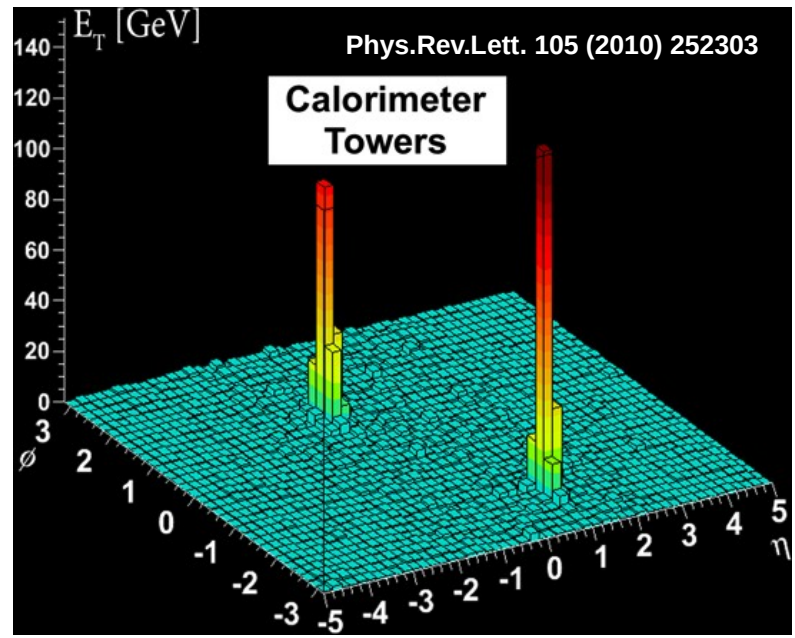
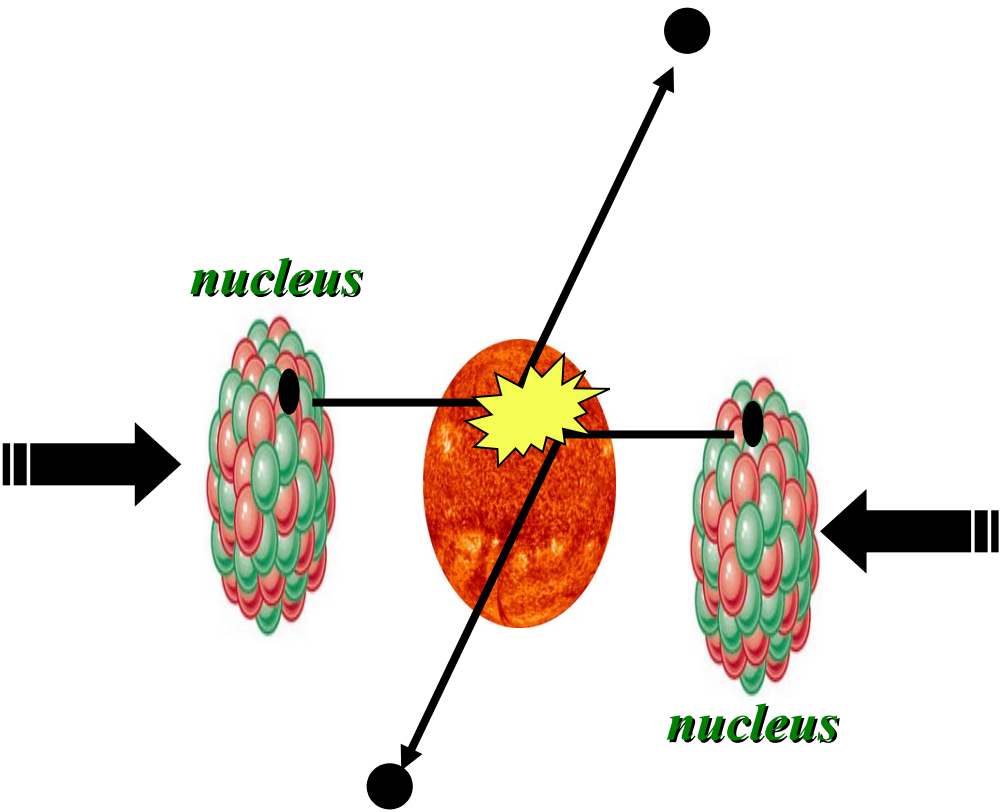
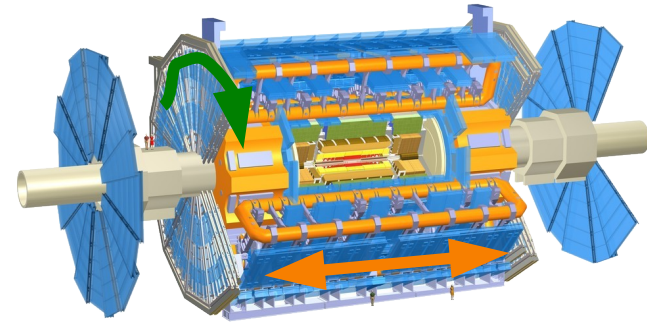
Probes of the Quark Gluon Plasma

Want a probe which traveled through the medium
QGP is short lived \rightarrow need a probe created in the collision
We expect the medium to be dense \rightarrow absorb/modify probe



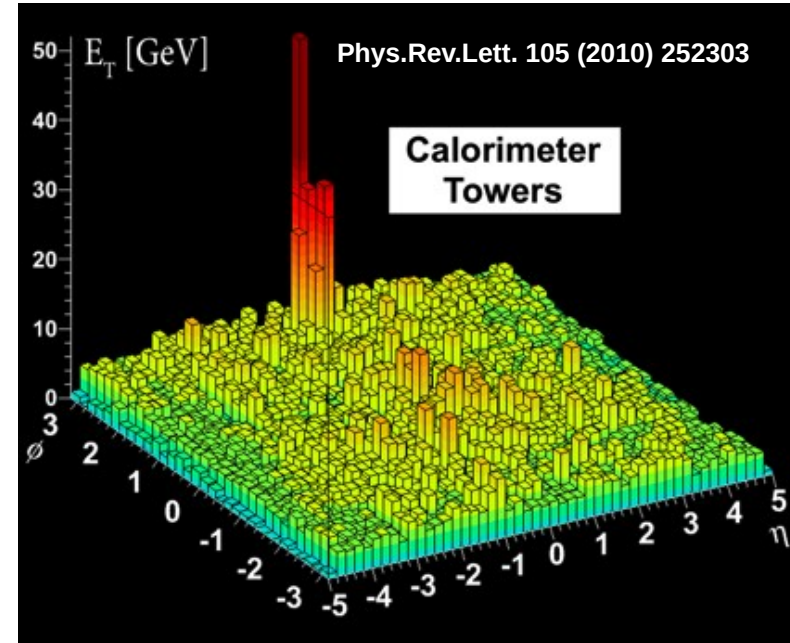
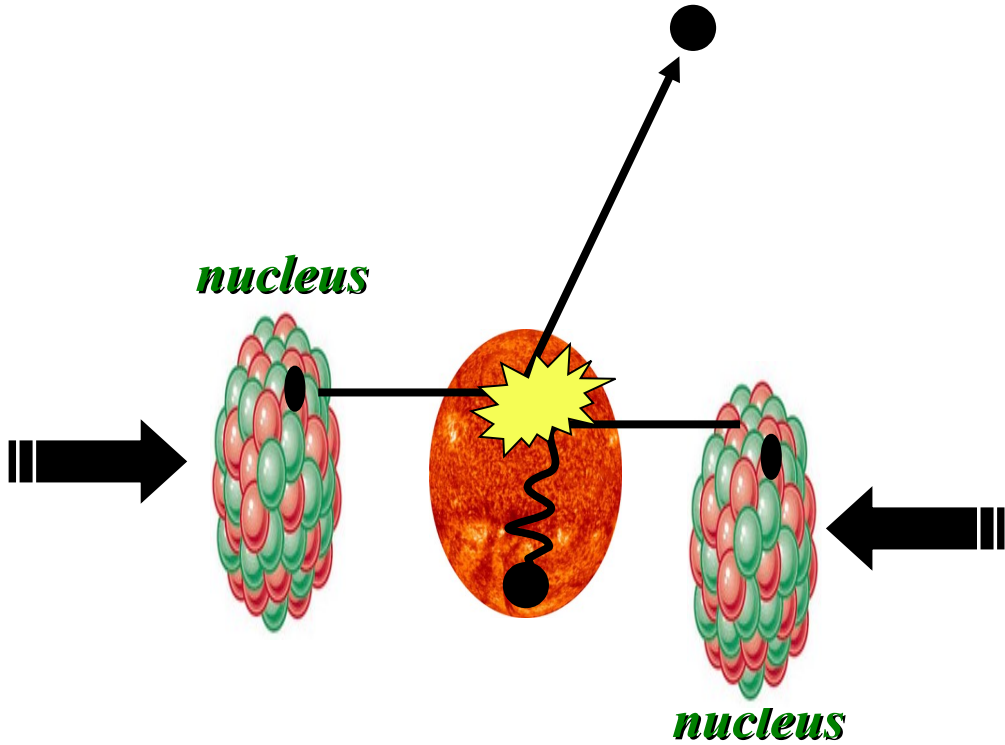
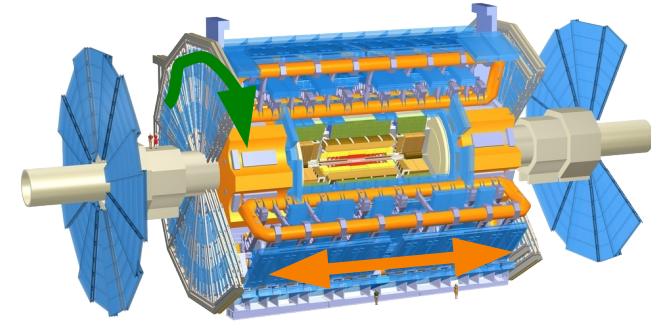
Probes of the Quark Gluon Plasma

ATLAS



Probes of the Quark Gluon Plasma

ATLAS

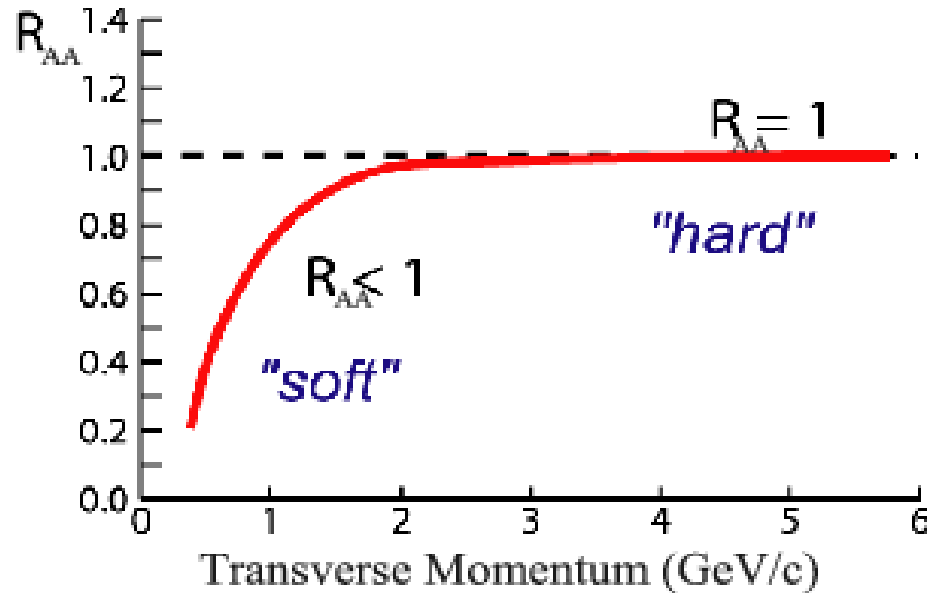


“Simple” example: Single hadrons

Nuclear modification factor

- Measure spectra of probe (jets) and compare to those in p+p collisions or peripheral A+A collisions
- If high- p_T probes (jets) are suppressed, this is evidence of jet quenching

$$R_{AA} = \frac{d^2 N_{AA}/dp_T d\eta}{T_{AA} d^2 \sigma^{pp}/dp_T d\eta}$$



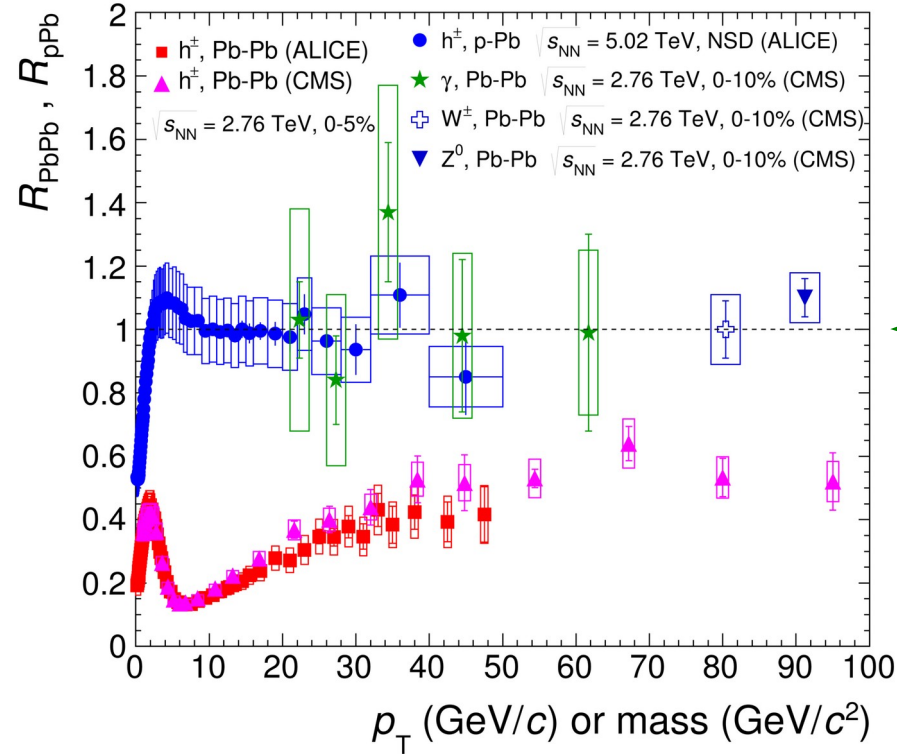
Enhancement

Suppression

Nuclear modification factor

Control →

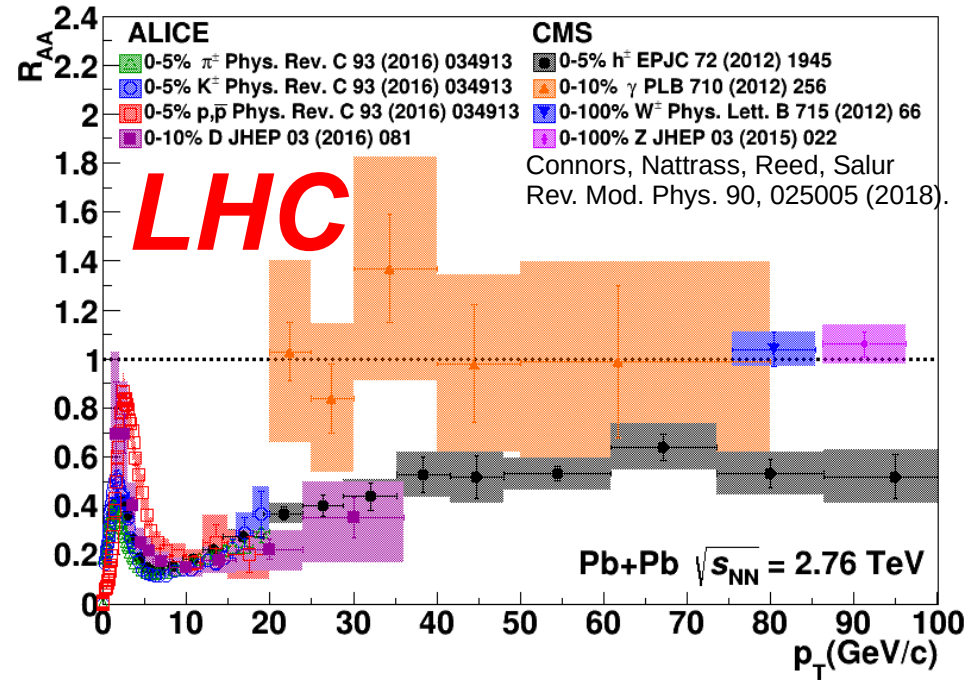
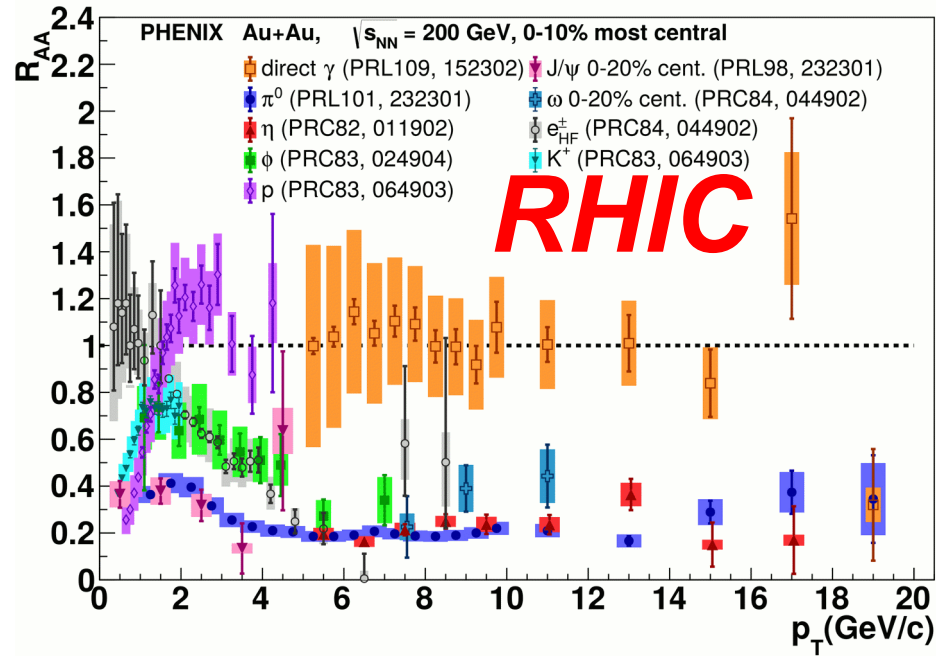
Probe →



← **Control**

- Charged hadrons (colored probes) suppressed in Pb—Pb
- Charged hadrons not suppressed in p—Pb at midrapidity
- Electroweak probes not suppressed in Pb—Pb

Nuclear modification factor R_{AA}



Electromagnetic probes – consistent with no modification – medium is transparent to them

Strong probes – significant suppression – medium is opaque to them
- even heavy quarks!

What am I measuring?
Definition of a jet

What is a jet?

What is a jet?

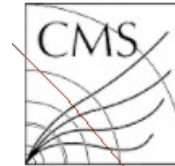
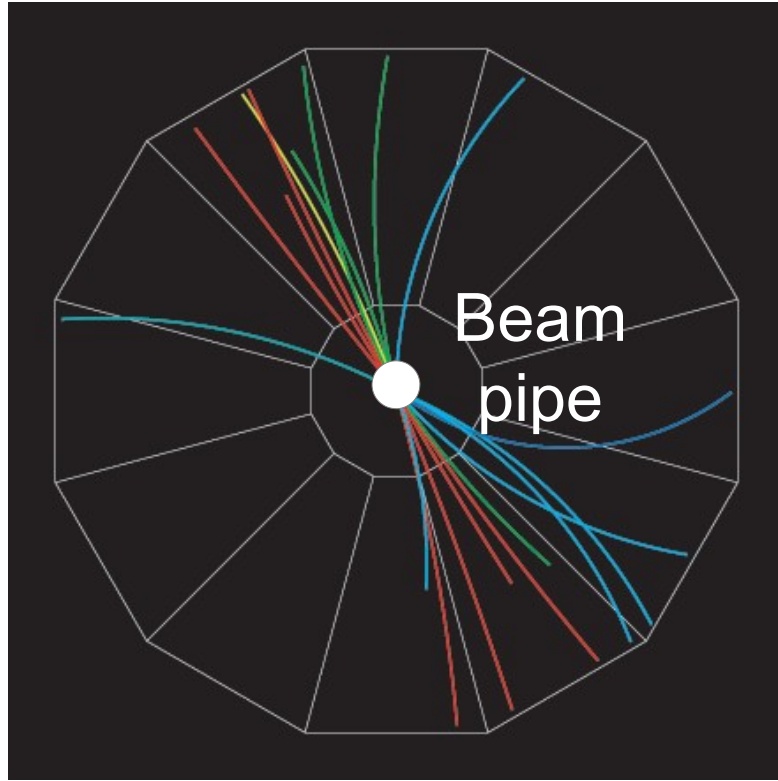
A measurement of a jet is a measurement of a parton.

What is a jet?

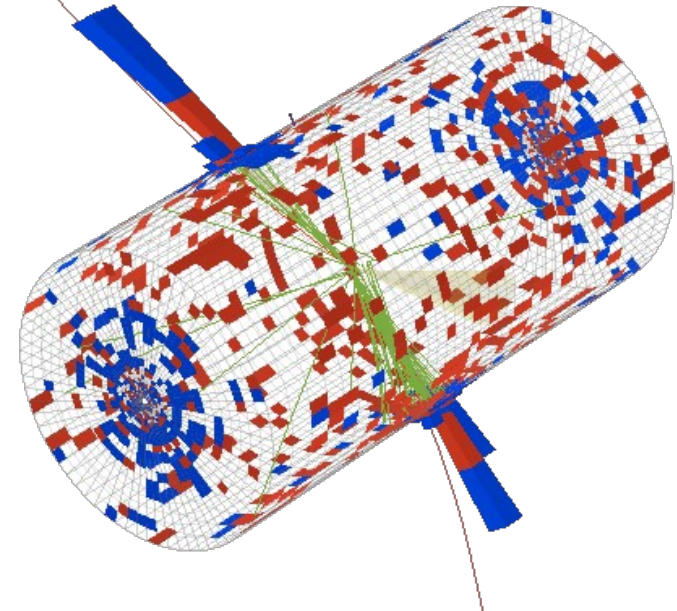
~~A measurement of a jet is a measurement of a
parton.~~

What is a jet?

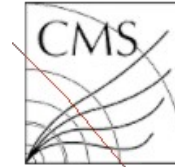
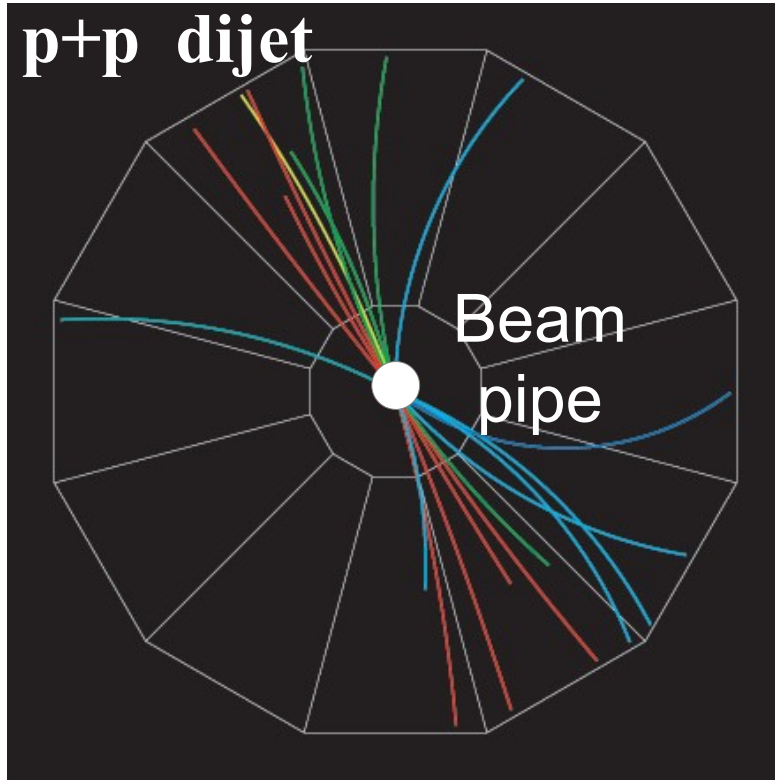
p+p dijet



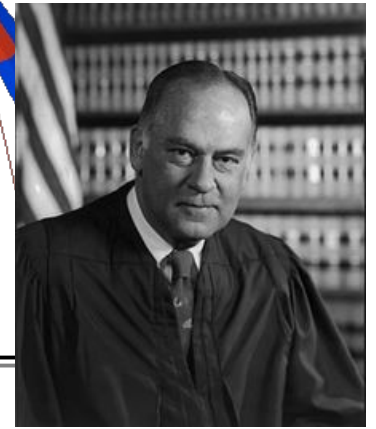
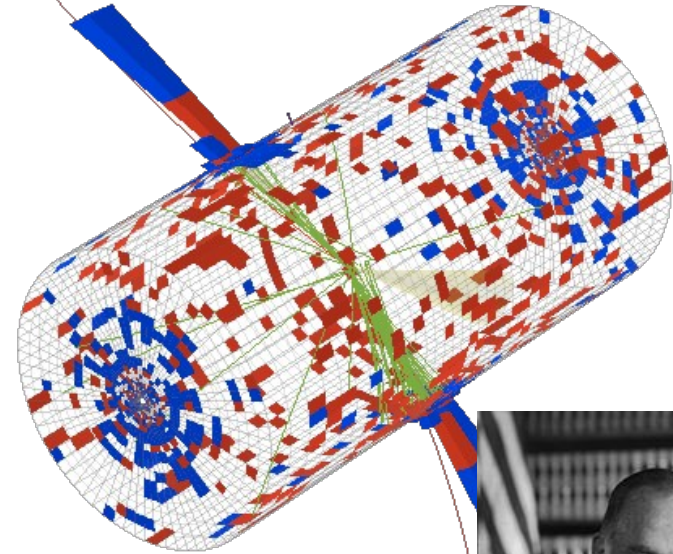
CMS Experiment at LHC, CERN
Data recorded: Fri Oct 5 12:29:33 2012 CEST
Run/Event: 204541 / 52508234
Lumi section: 32



What is a jet?



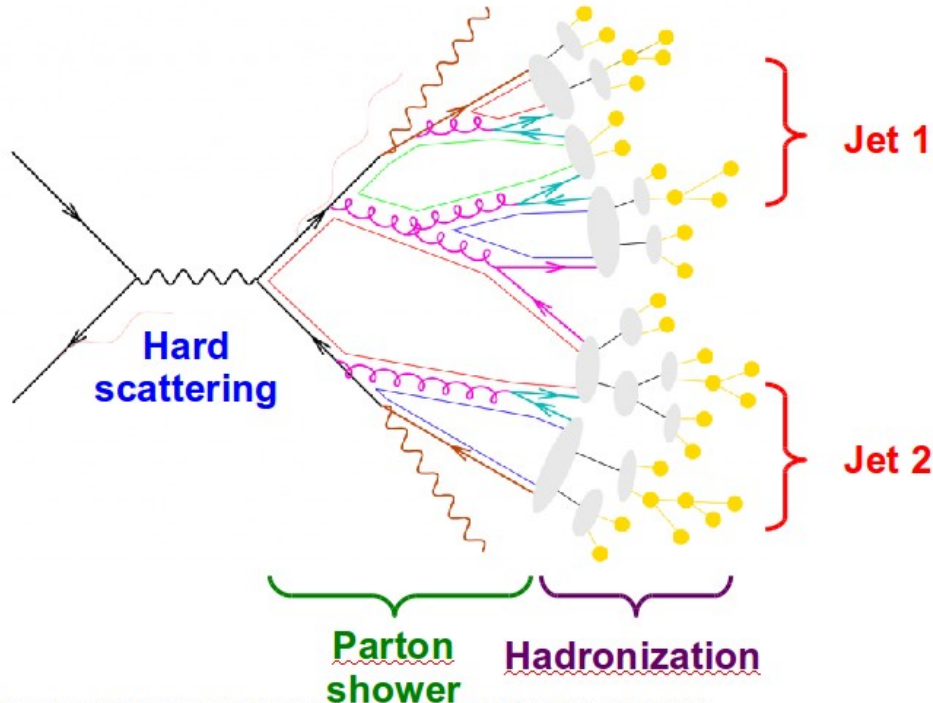
CMS Experiment at LHC, CERN
Data recorded: Fri Oct 5 12:29:33 2012 CEST
Run/Event: 204541 / 52508234
Lumi section: 32



“I know it when I see it”

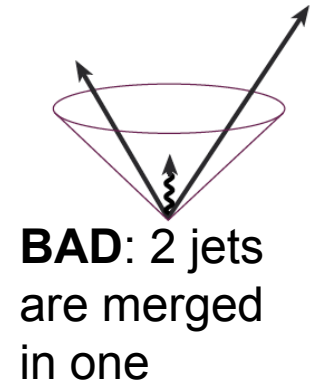
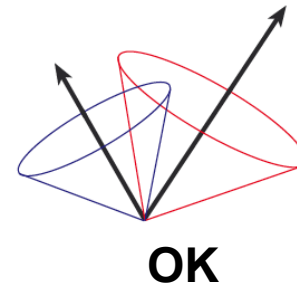
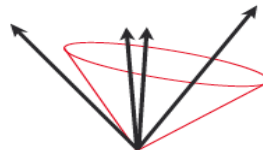
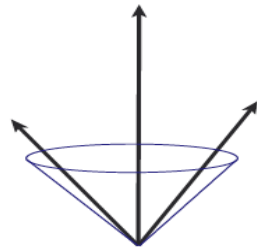
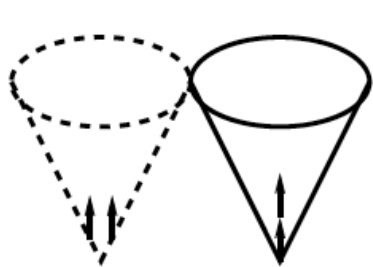
US Supreme Court Justice Potter Stewart, *Jacobellis v. Ohio*

Jets in principle



- Jet measures **partons**
- Hadronic degrees of freedom are integrated out
- Algorithms are infrared and collinear safe

Image from <http://www.gk-eichtheorien.physik.uni-mainz.de/Dateien/Zepfenfeld-3.pdf>



Jet finding *in pp collisions*

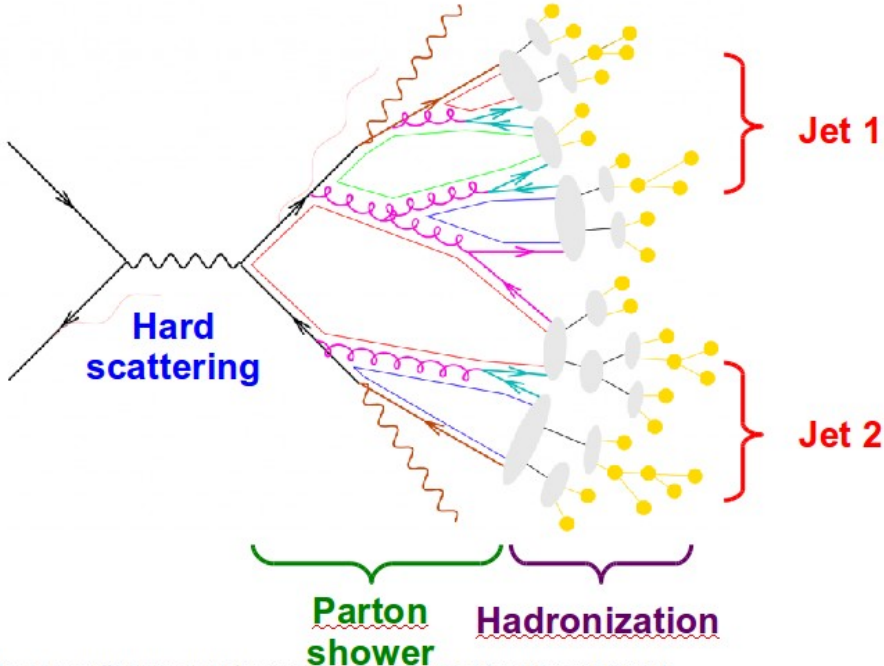
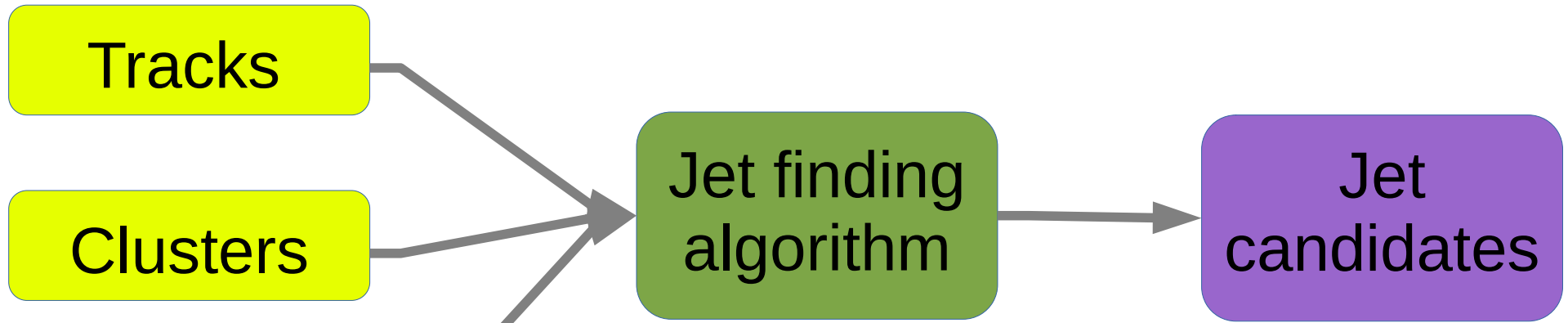


Image from <http://www.gk-eichtheorien.physik.uni-mainz.de/Dateien/Zeppenfeld-3.pdf>

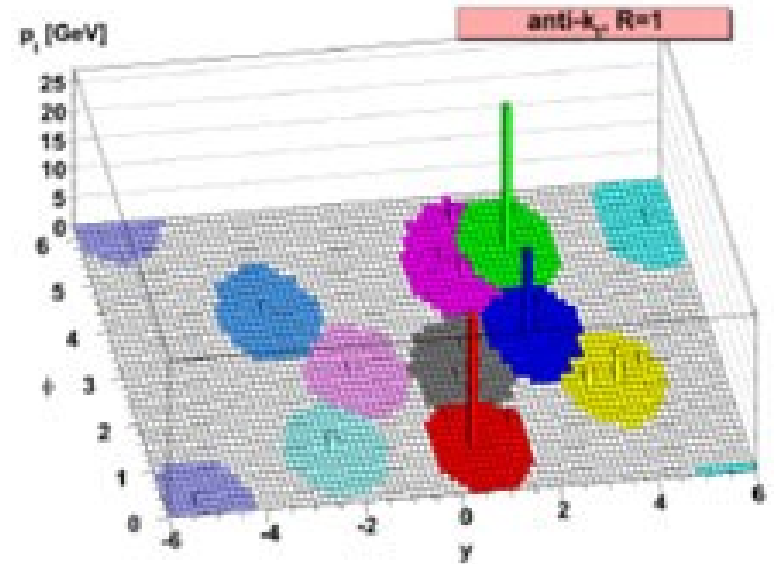
- Jet finder: groups final state particles into jet candidates
 - Anti- k_T algorithm
[JHEP 0804 \(2008\) 063 \[arXiv:0802.1189\]](#)
- Depends on hadronization
- Ideally
 - Infrared safe
 - Collinear safe

Snowmass Accord: Theoretical calculations and experimental measurements should use the same jet finding algorithm. Otherwise they will not be comparable.

Jet finding algorithms



- Any list of objects works as input
- Use the same algorithm on theory & experiment
- Output only as good as input



k_T jet finding algorithm

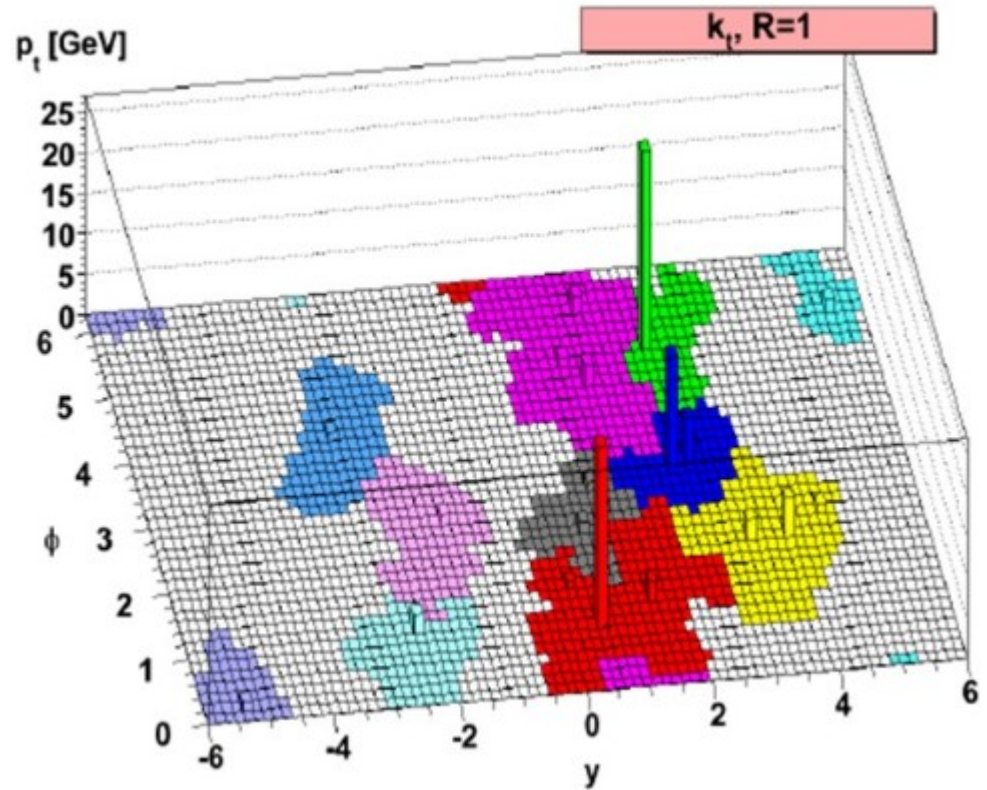
Particles, clusters

k_T algorithm

$$k_T = p_T, \Delta R_{ij} = \sqrt{(\eta_i - \eta_j)^2 + (\phi_i - \phi_j)^2}$$

- For all i, j calculate:
 $d_{ij} = \min(p_{T,i}^2, p_{T,j}^2) \frac{\Delta R_{ij}^2}{R^2}$
 - $d_{iB} = p_{T,i}^2$
 - Combine smallest d_{ij} .
If d_{iB} smallest, $d_{iB} \rightarrow$ jet
- Repeat until no particles left

Jet candidates



anti- k_T jet finding algorithm

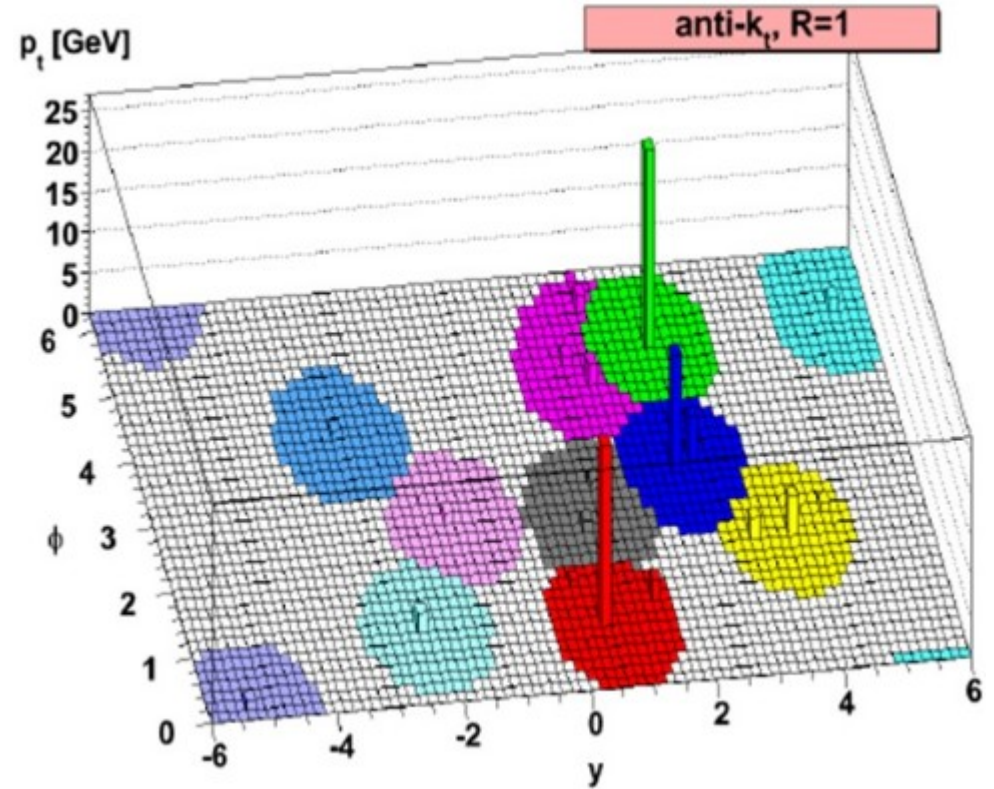
Particles, clusters

k_T algorithm

$$k_T = p_T, \Delta R_{ij} = \sqrt{(\eta_i - \eta_j)^2 + (\phi_i - \phi_j)^2}$$

- For all i, j calculate:
$$d_{ij} = \min(p_{T,i}^{-2}, p_{T,j}^{-2}) \frac{\Delta R_{ij}^2}{R^2}$$
 - $d_{iB} = p_{T,i}^{-2}$
 - Combine smallest d_{ij} .
If d_{iB} smallest, $d_{iB} \rightarrow$ jet
- Repeat until no particles left

Jet candidates



Cambridge/Aachen jet finding algorithm

Particles, clusters

k_T algorithm

$$k_T = p_T, \Delta R_{ij} = \sqrt{(\eta_i - \eta_j)^2 + (\phi_i - \phi_j)^2}$$

- For all i, j calculate:

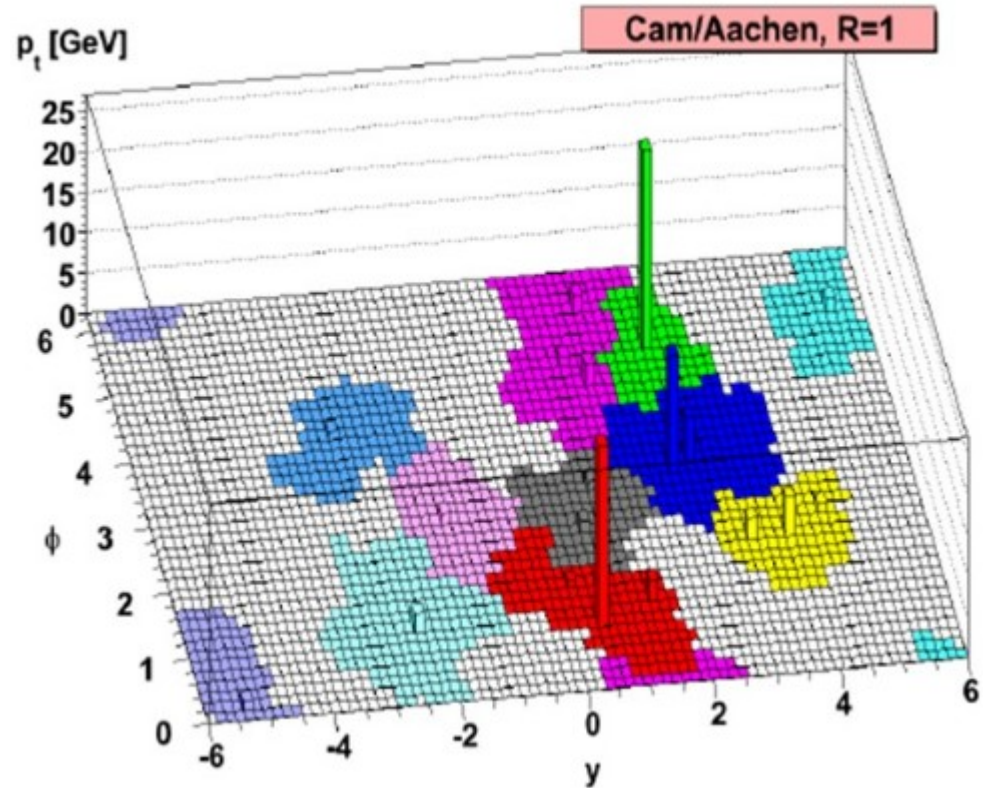
$$d_{iB} = 1 \quad d_{ij} = \frac{\Delta R_{ij}^2}{R^2}$$

- Combine smallest d_{ij} .

If d_{iB} smallest, $d_{iB} \rightarrow$ jet

Repeat until no particles left

Jet candidates



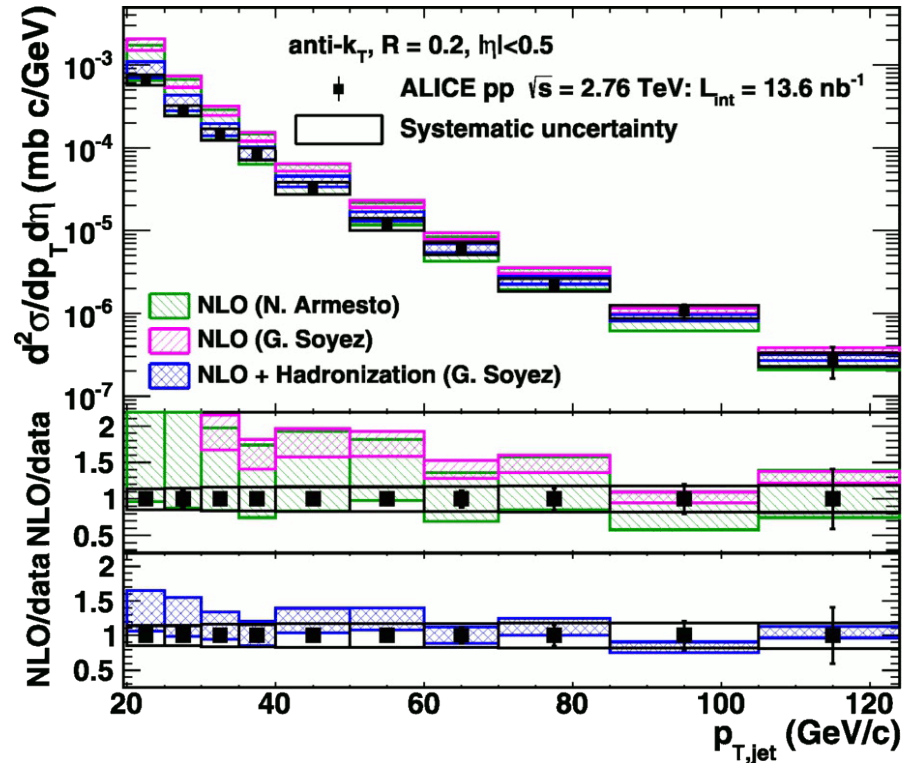
A jet is what a jet finder finds.

Jet cross-section in pp

$\sqrt{s} = 2.76$ TeV, $R = 0.2$ Inclusive

arXiv:1301.3475

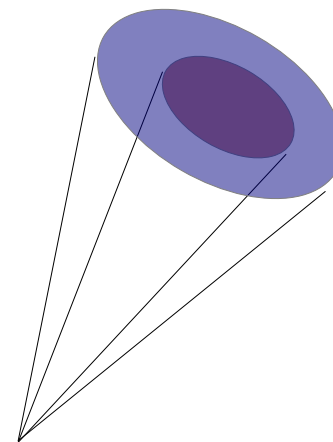
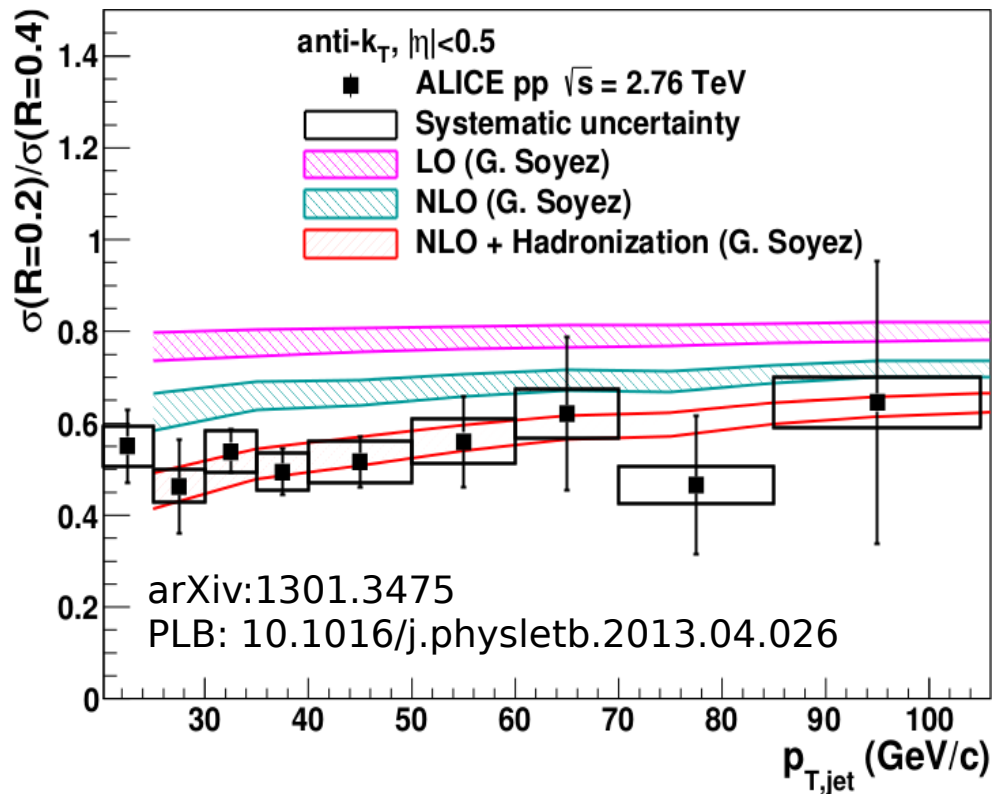
PLB: 10.1016/j.physletb.2013.04.026



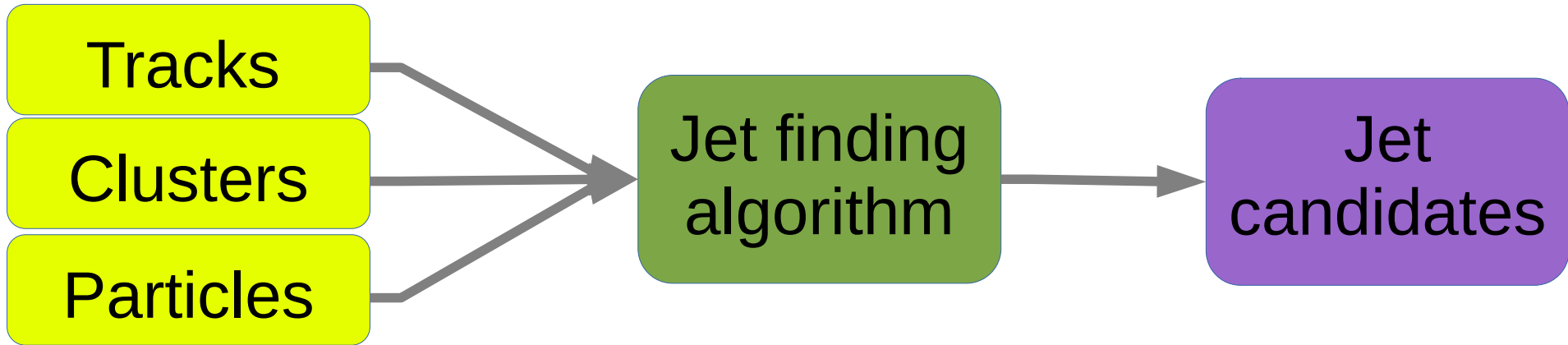
- Green and magenta bands: NLO on Parton level
- Blue band: NLO + hadronization
- Hadronization calculations necessary to describe data

Jet ratios in pp

$\sqrt{s} = 2.76$ TeV, $R = 0.2, 0.4$ Inclusive



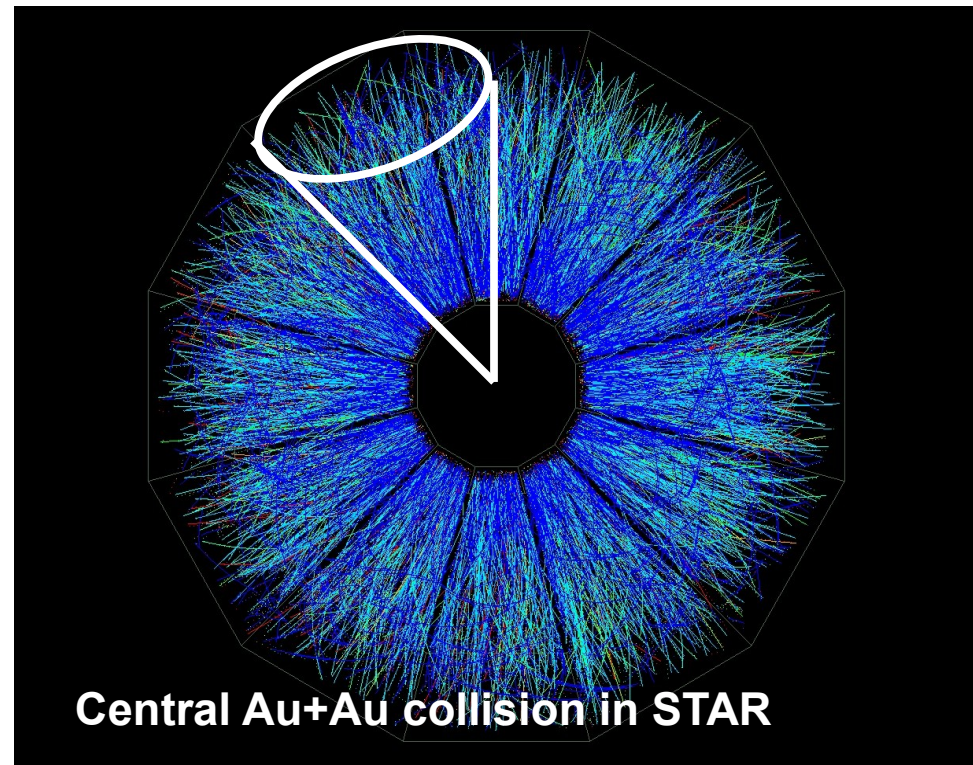
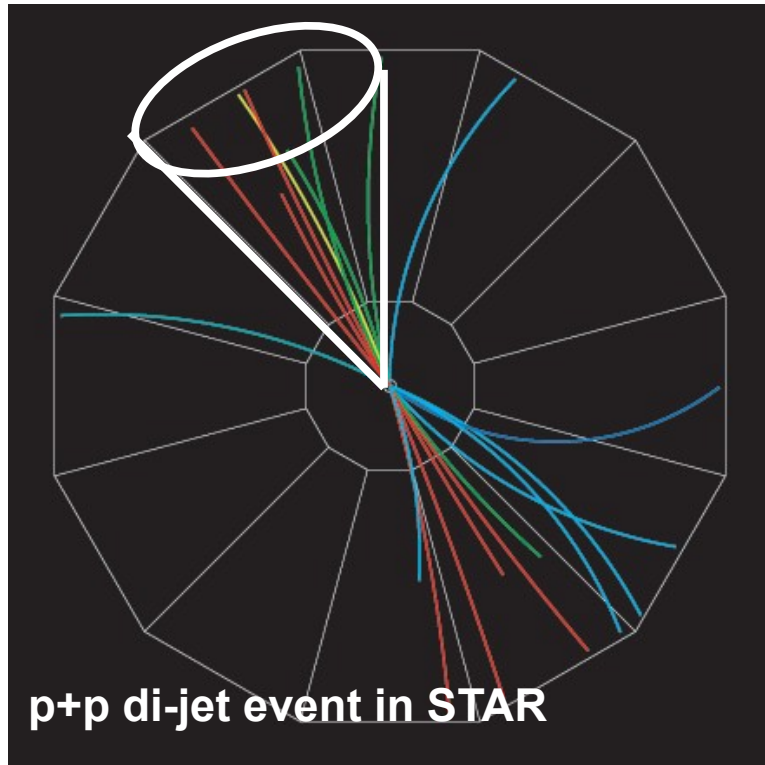
Mini-summary



- Jets are not partons
- Good jet finders:
 - Infrared and colinear safe
 - k_T , anti- k_T , Cambridge/Aachen, SIScone
- Jet is defined by jet finder, its parameters
- PDFs, fragmentation functions non-perturbative
→ all jet measurements sensitive to somewhat non-perturbative effects
- Good agreement between theory and experiment

Jets in A+A collisions
What assumptions am I making?

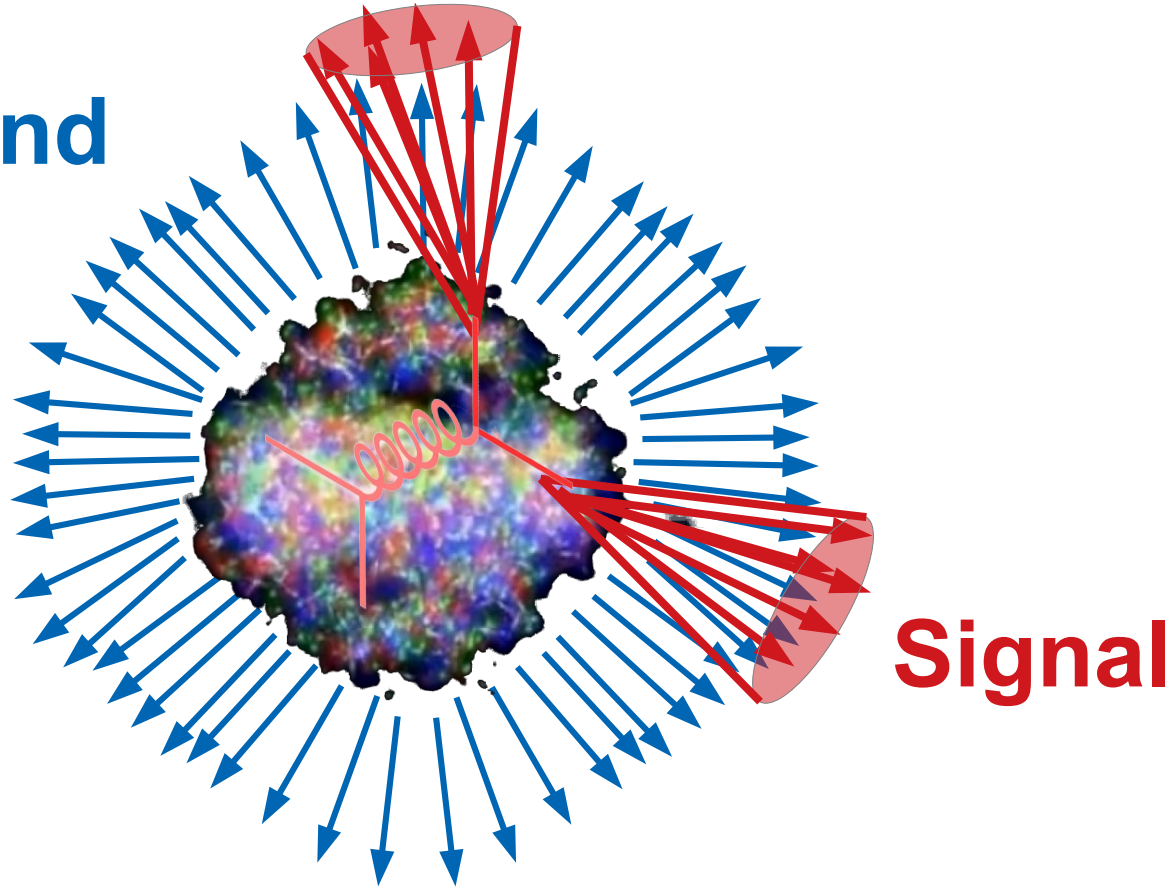
$p+p$ vs $A+A$



Signal vs Background:

The standard paradigm

Background

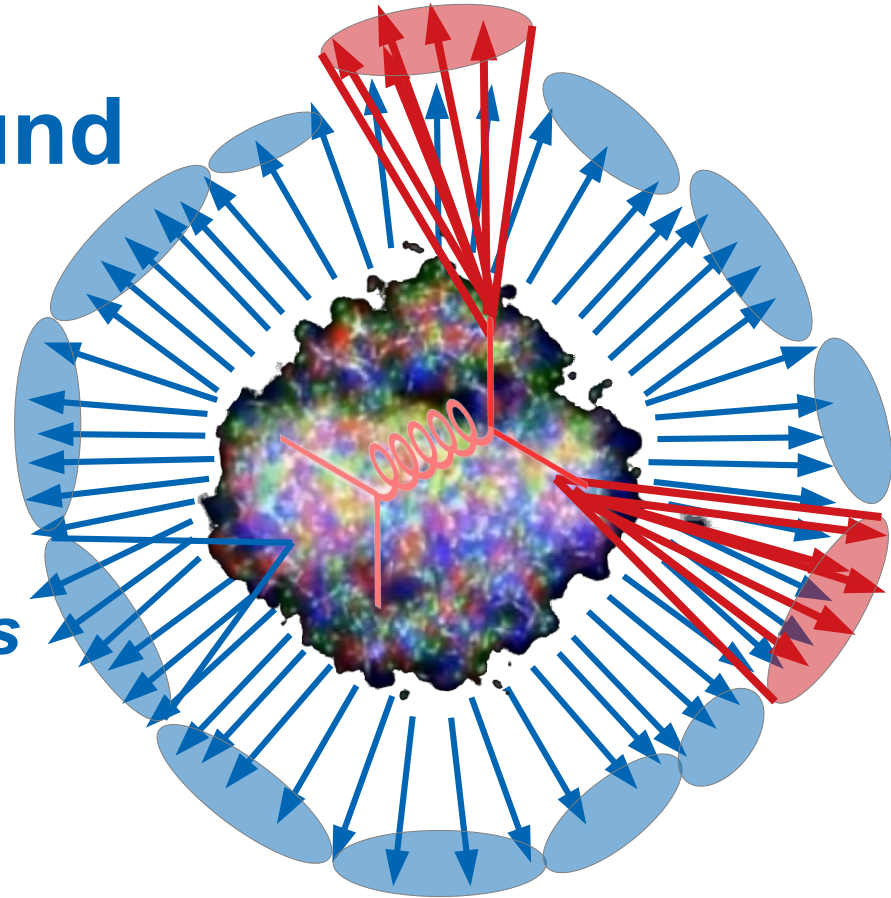


Signal vs Background:

The standard paradigm

Background

Combinatorial jets



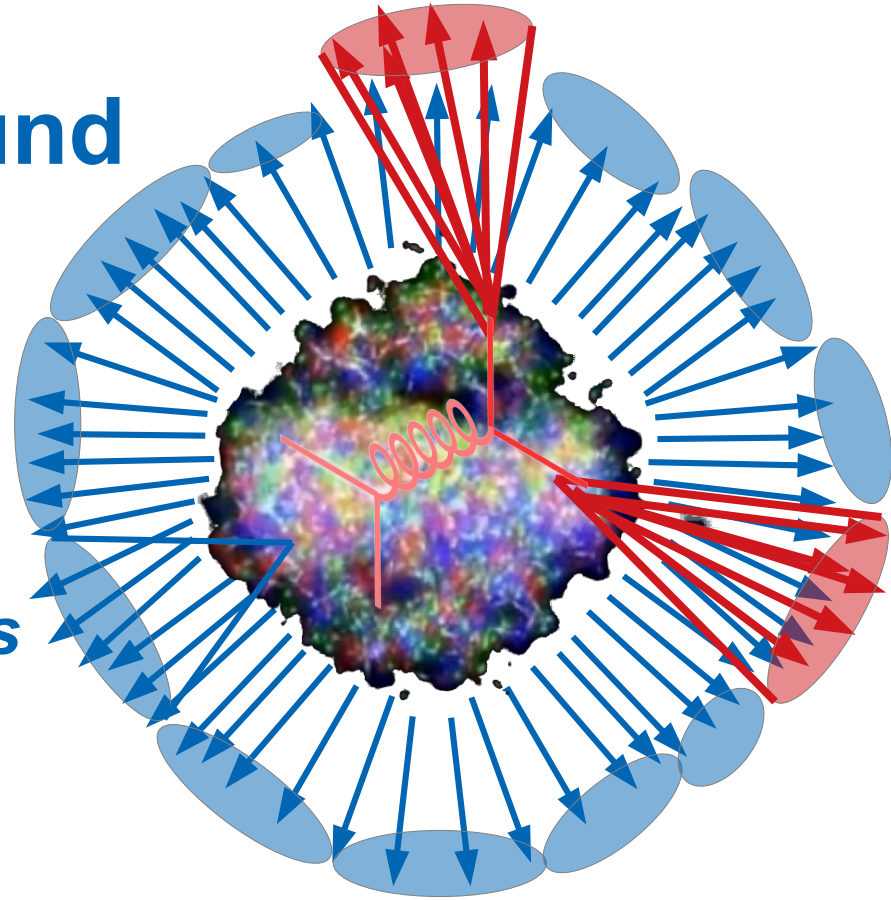
Signal

Signal vs Background:

The standard paradigm

Background

**Combinatorial jets
= “fake” jets**



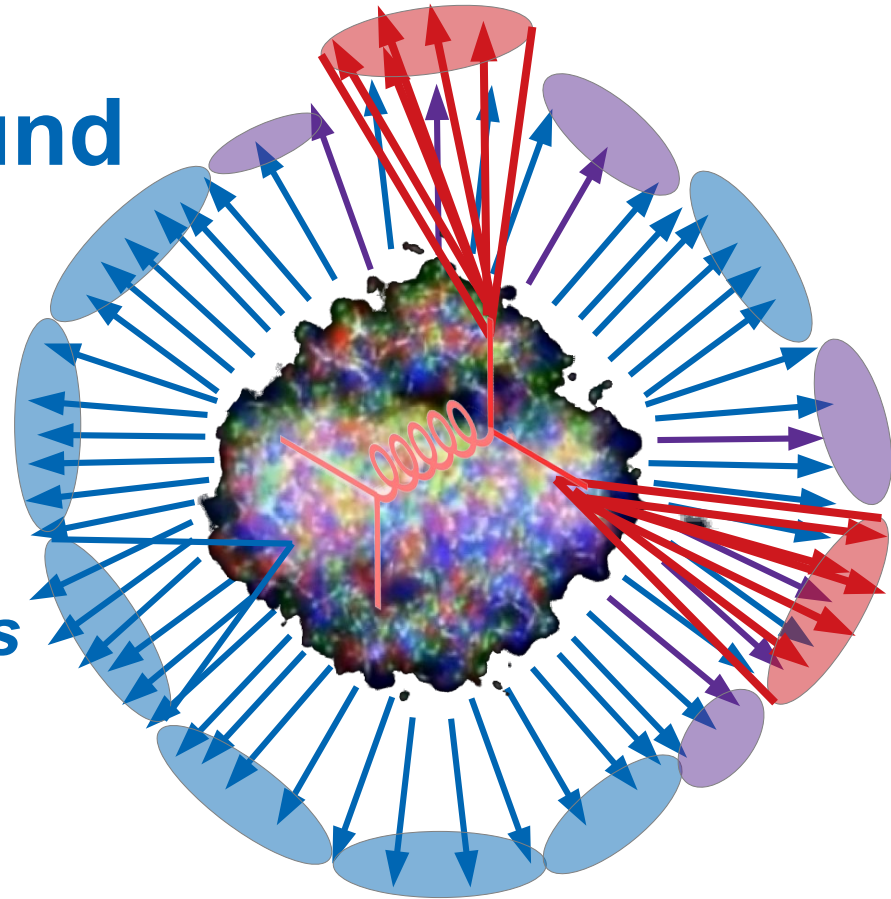
Signal

Signal vs Background:

The standard paradigm

Background

Combinatorial jets



Signal

***Some gray areas**

Jet finding in AA collisions

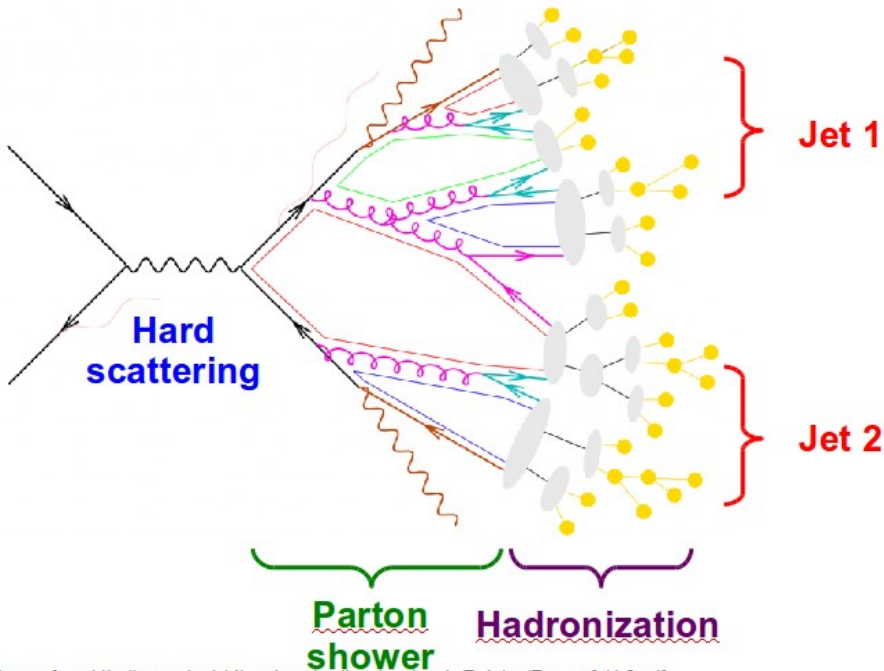
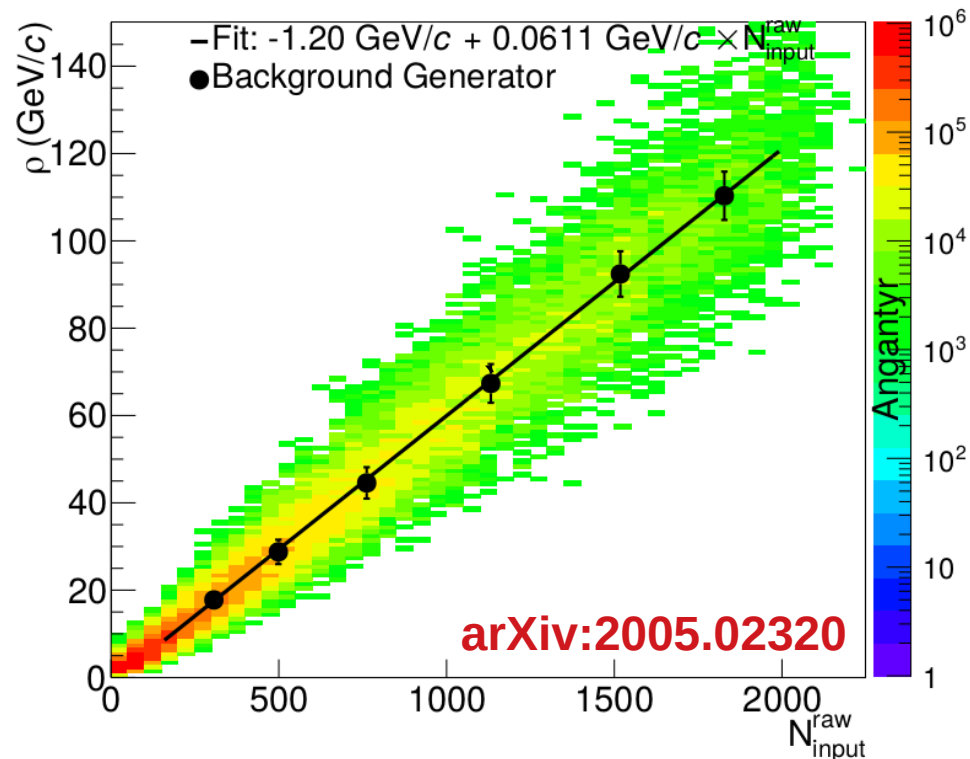
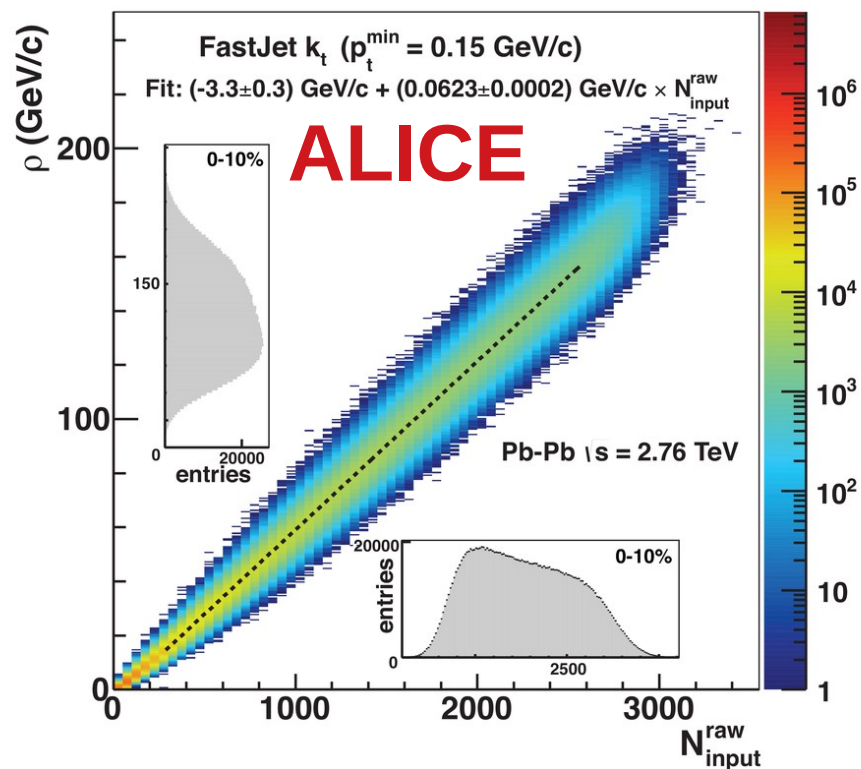


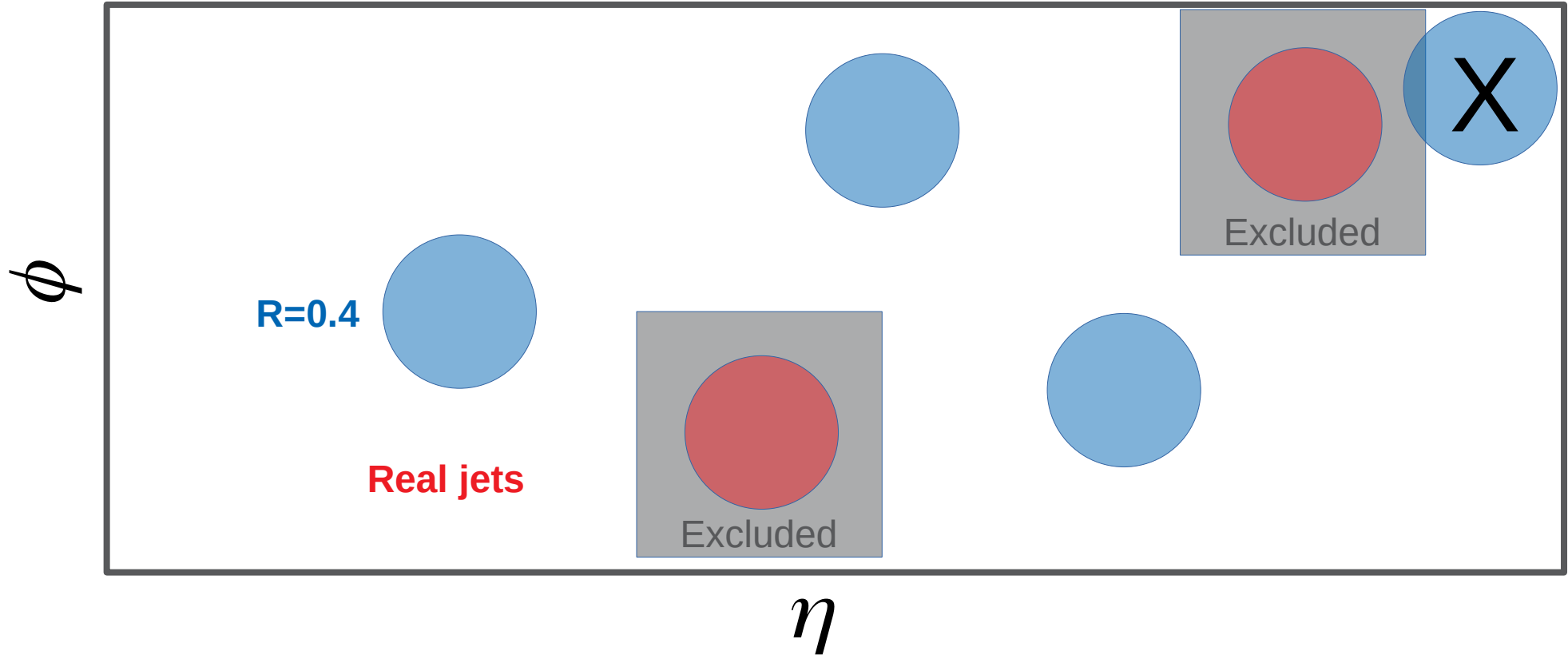
Image from <http://www.gk-eichtheorien.physik.uni-mainz.de/Dateien/Zeppenfeld-3.pdf>

- Jet finder: groups final state particles into jet candidates
 - Anti- k_T algorithm
JHEP 0804 (2008) 063 [arXiv:0802.1189]
- Combinatorial jet candidates
- Energy smearing from background
- Sensitive to methods to suppress combinatorial jets and correct energy
- Focus on narrow/high energy jets

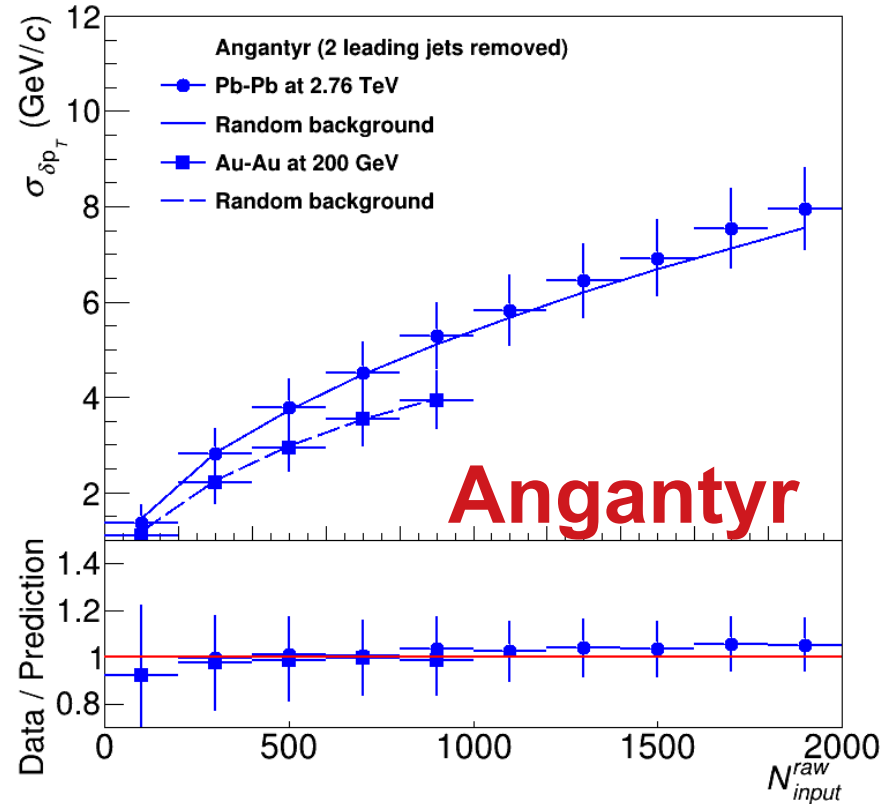
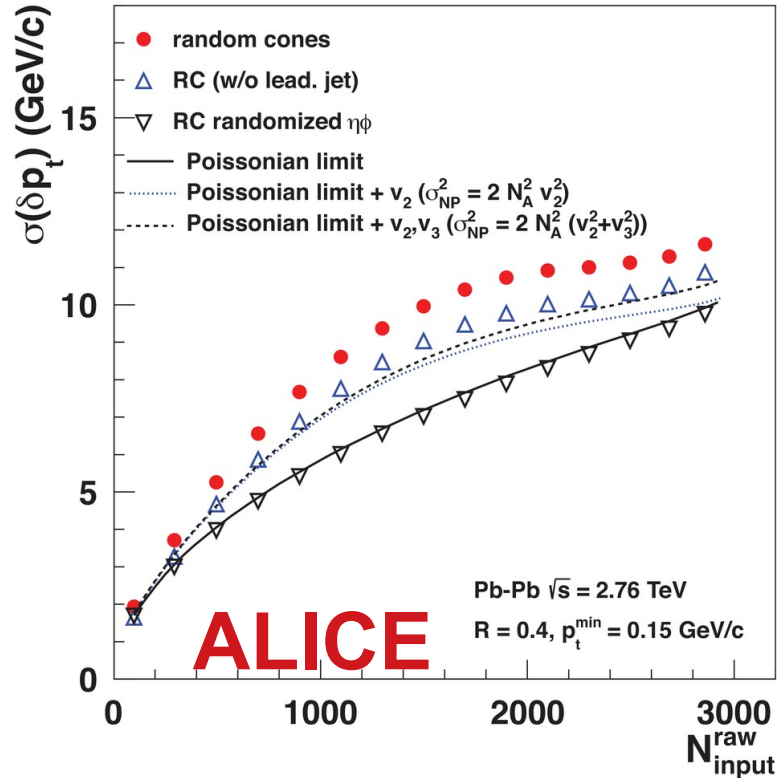
Background density ρ



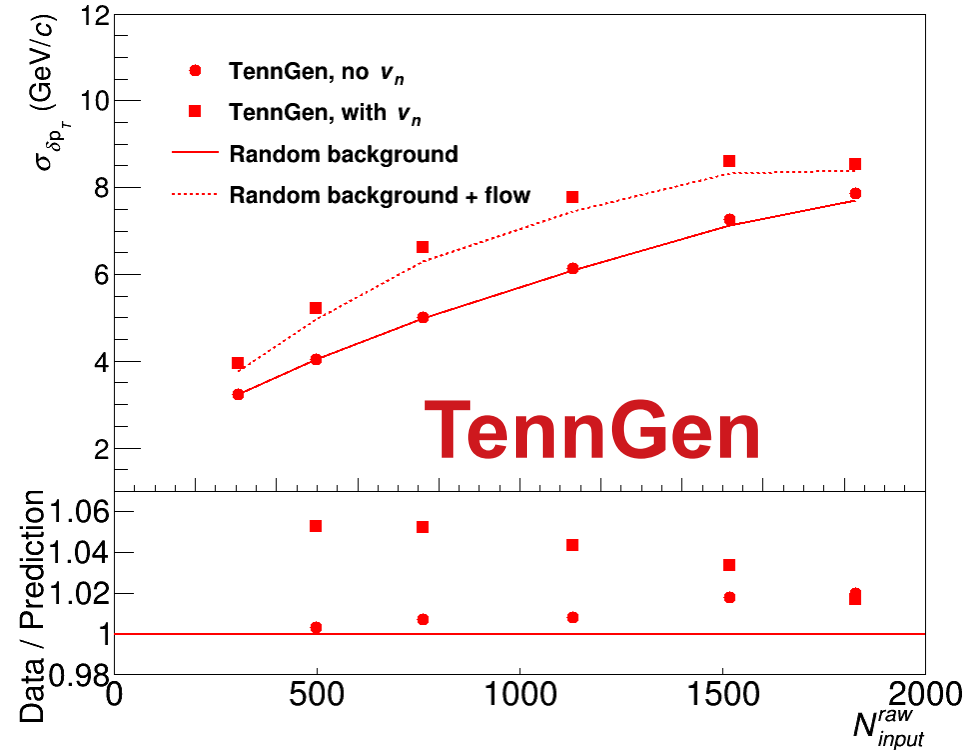
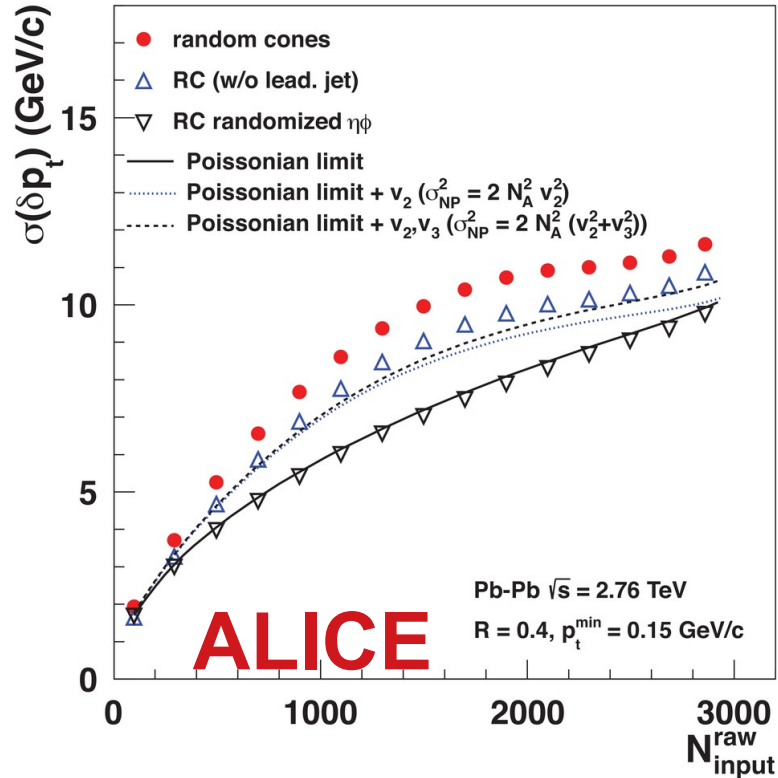
Random cones



Width vs multiplicity

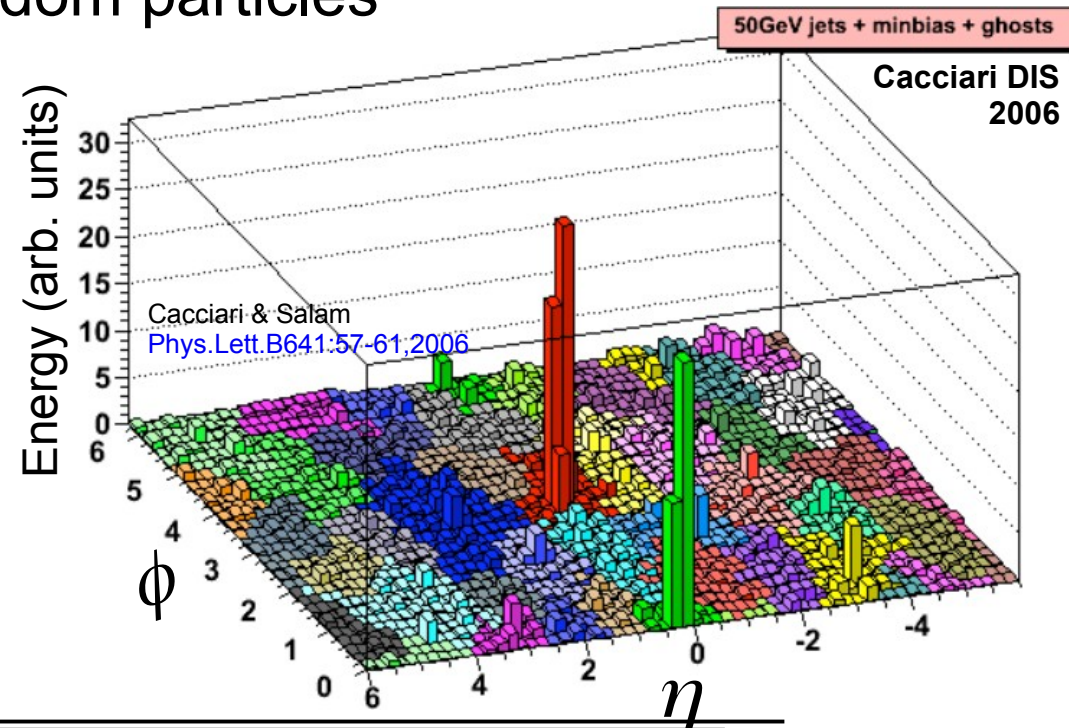


Width vs multiplicity



Mini-summary

- Jet finders put all input clusters, tracks in a jet candidate
- Background is *dominated* by random particles
 - But 5% effects from flow
- Models have background too!
 - Sensitive to multiplicity, shape of spectrum

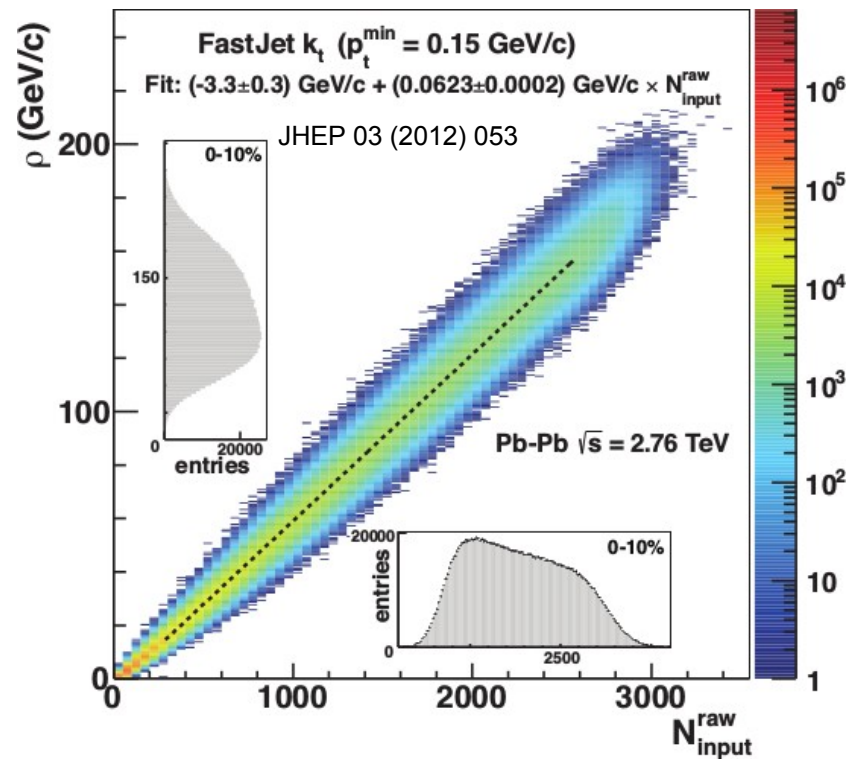


Jets in A+A collisions: Dealing with background

Focus on smaller angles

- Pros
 - Background is smaller
 - Background fluctuations smaller
- Cons:
 - Modifications expected at higher R
 - Biases sample towards quarks

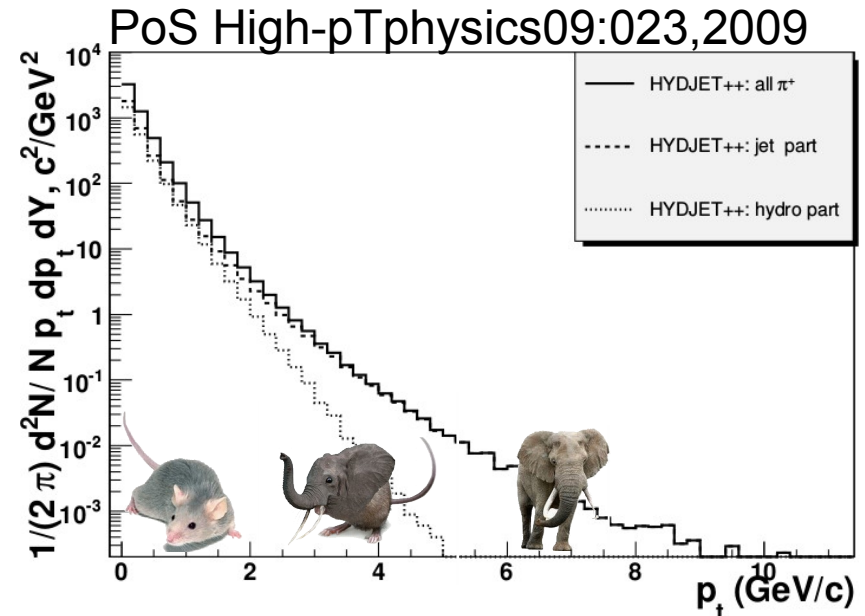
Aside: “quark” and “gluon” jet only defined at leading order.



Focus on high p_T

- Pros:
 - Reduces combinatorial background
- Cons:
 - Cuts signal where we expect modifications
 - Could bias towards partons which have not interacted
 - Biases sample towards quark jets

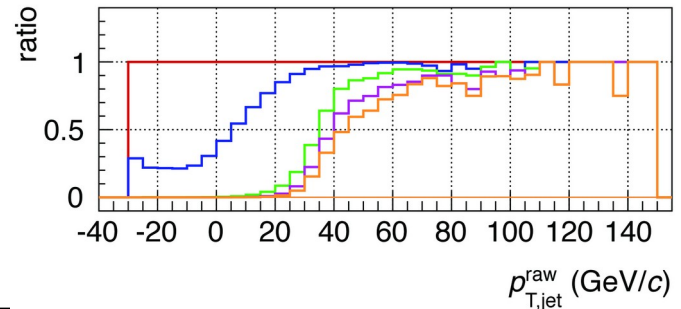
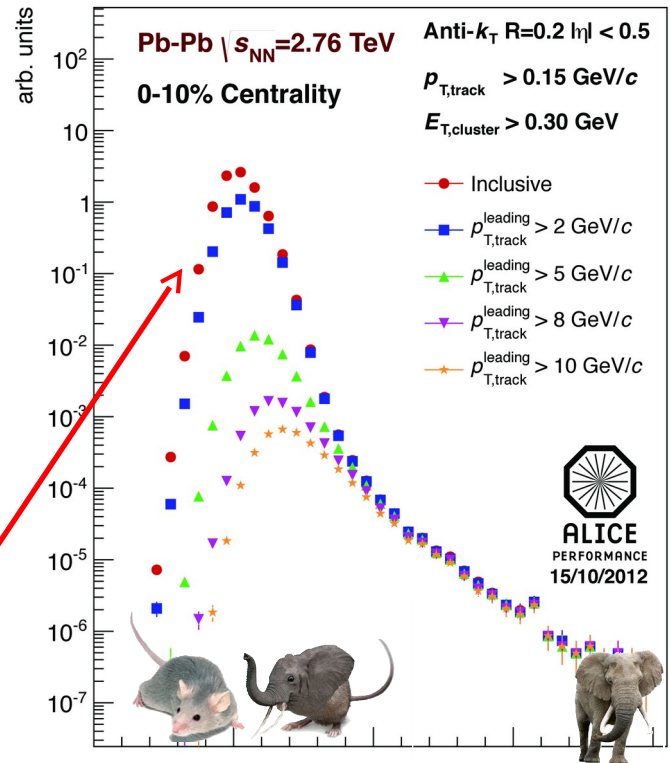
“Quark” and “gluon” jets only defined at leading order!



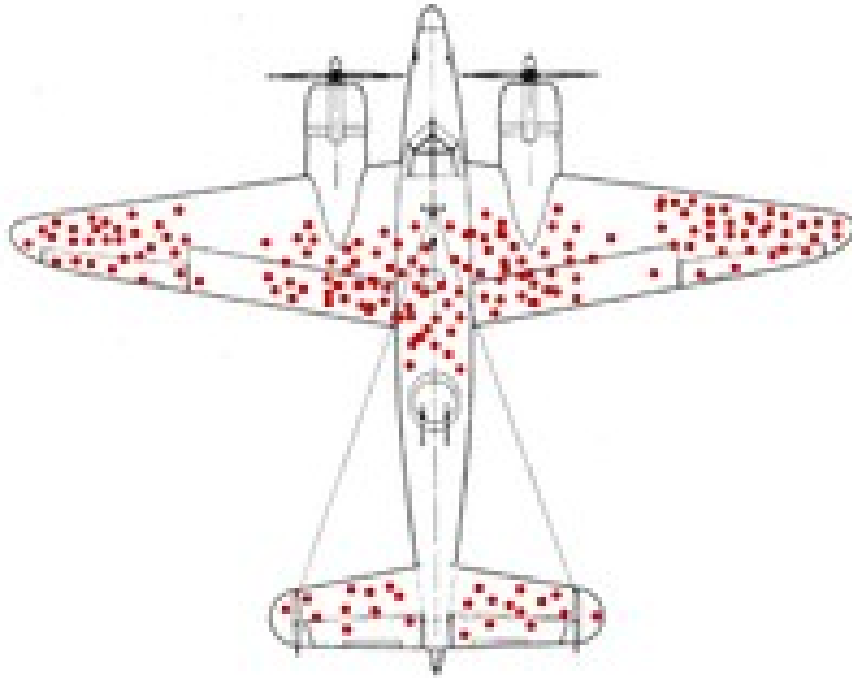
Area-based subtraction

- ALICE/STAR
- Require leading track $p_T > 5 \text{ GeV}/c$
 - Suppresses combinatorial “jets”
 - Biases fragmentation
- No threshold on constituents
- Limited to small R

Combinatorial “jets”



Survivor bias



- **WWII Example:** holes planes returning indicate where it's *safer* to get hit
- We're looking at the jets which *remain*



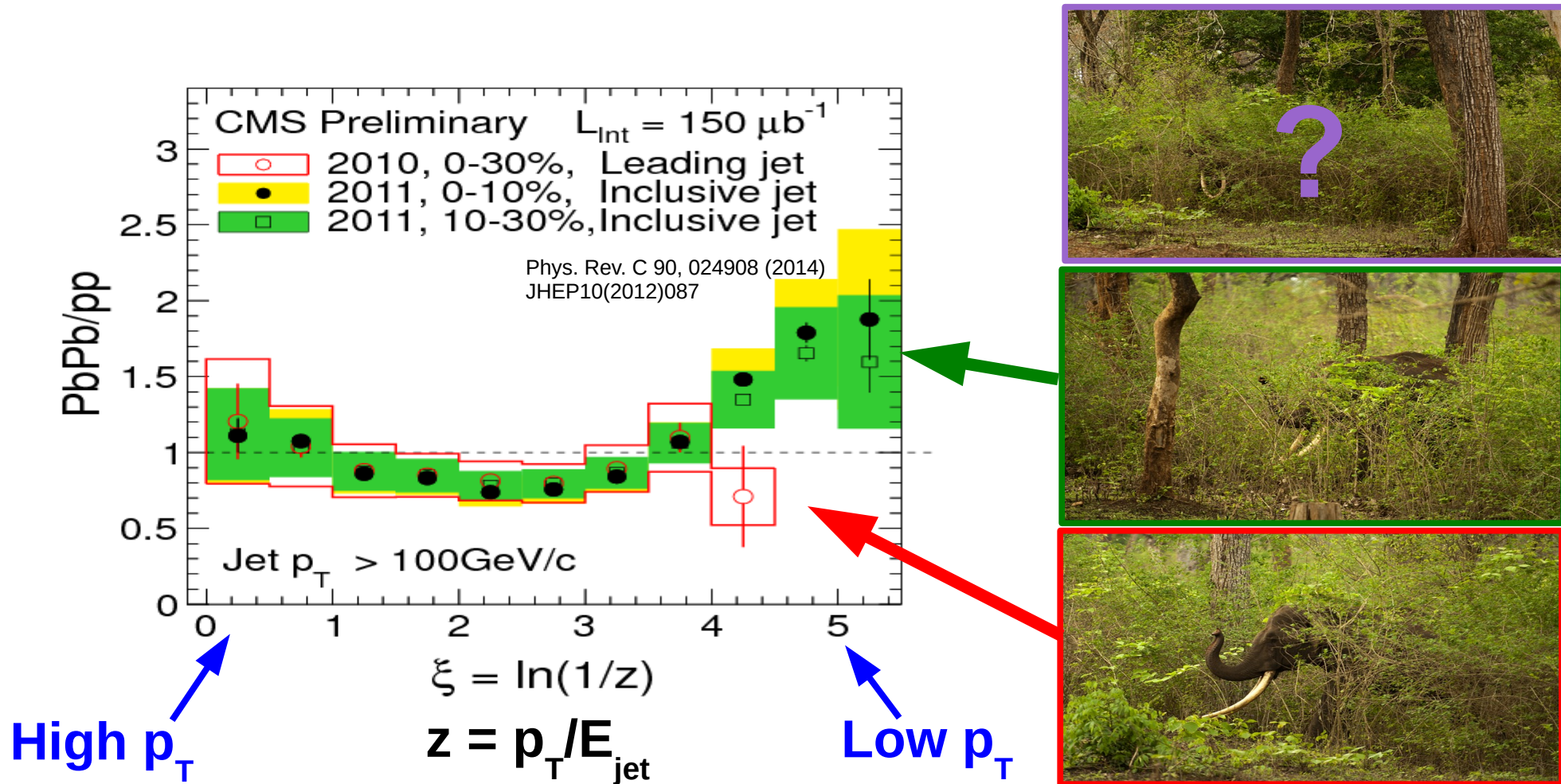


**What you see depends on what you're
looking for**

Bias & background

- **Experimental background subtraction methods:** complex, make assumptions, apply biases
- **Survivor bias:** Modified jets probably look more like the medium
- **Quark/Gluon bias:**
 - Quark jets are narrower, have fewer tracks, fragment harder [Z Phys C 68, 179-201 (1995), Z Phys C 70, 179-196 (1996),]
 - Gluon jets reconstructed with k_T algorithm have more particles than jets reconstructed with anti- k_T algorithm [Phys. Rev. D 45, 1448 (1992)]
 - Gluon jets fragment into more baryons [EPJC 8, 241-254, 1998]
- **Fragmentation bias:** Experimental measurements explicitly select jets with hard fragments

What you see depends on where you look

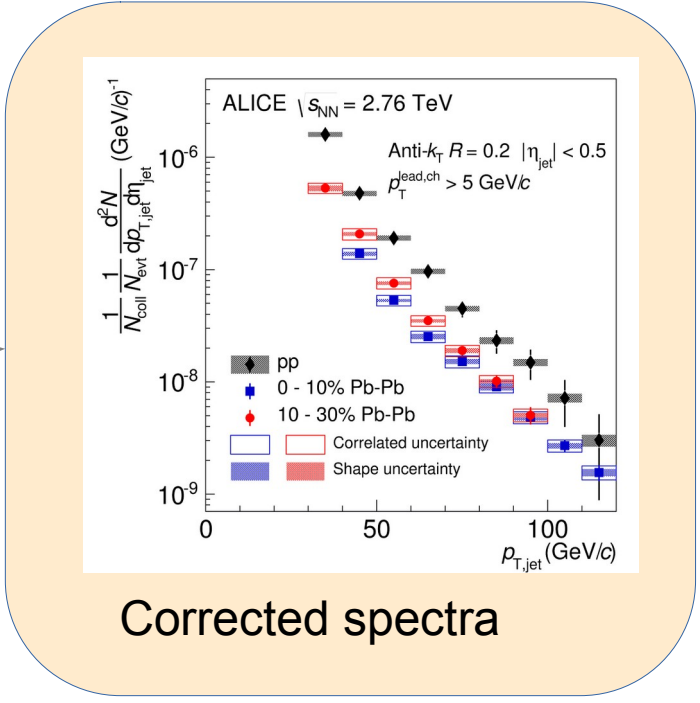
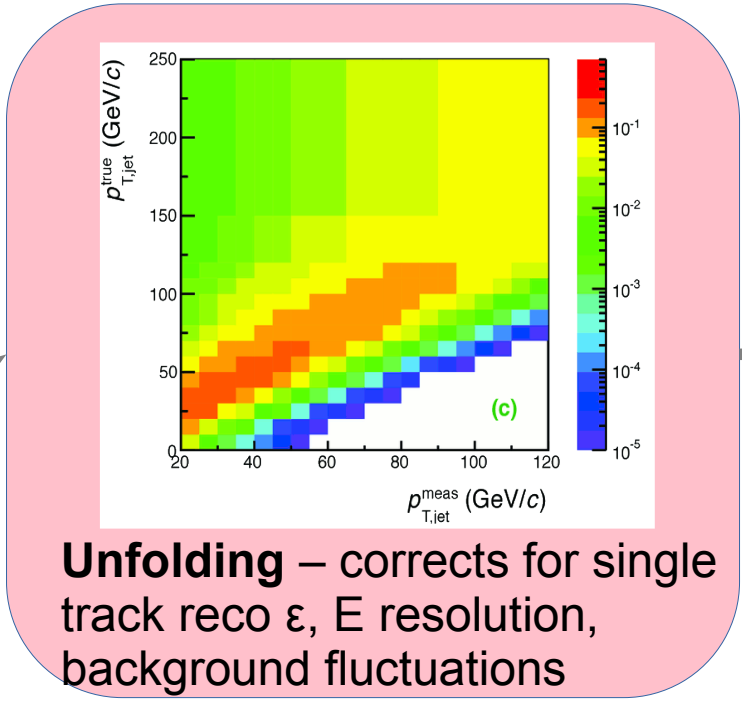
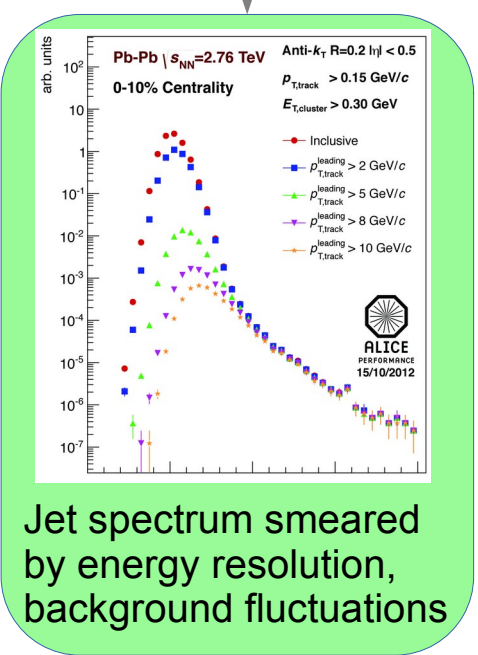
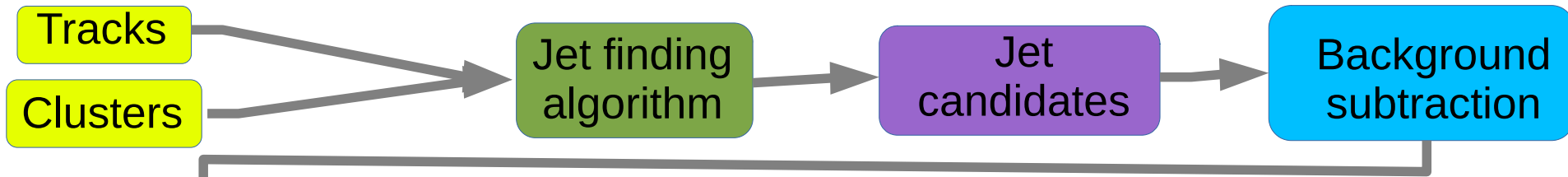


Mini-summary

- Most studies do one or more of the following:
 - Explicitly apply a (non-perturbative) bias
 - *Implicitly* apply a (non-perturbative) bias
 - Focus on small R
 - Focus on high pT
- May also → survivor bias
- Background subtraction should be part of definition of algorithm

What are the dominant uncertainties?

Analysis steps



Unfolding

$$\vec{v} = R\vec{\mu} + \vec{\beta}$$

$$v_i = \sum_{j=1}^M (R_{ij}\mu_j) + \beta_i$$

- $\vec{\mu}$: the “true” histogram
- \vec{v} : the actual data we measure
- $\vec{\beta}$: background
- R : the response matrix

Simple Solution (Inversion)

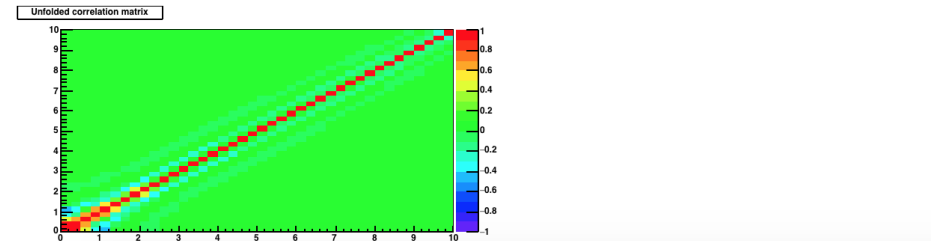
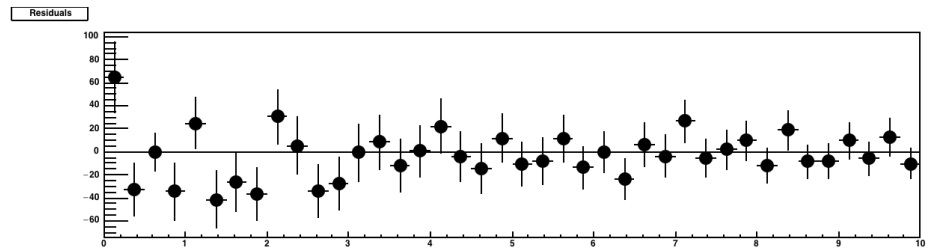
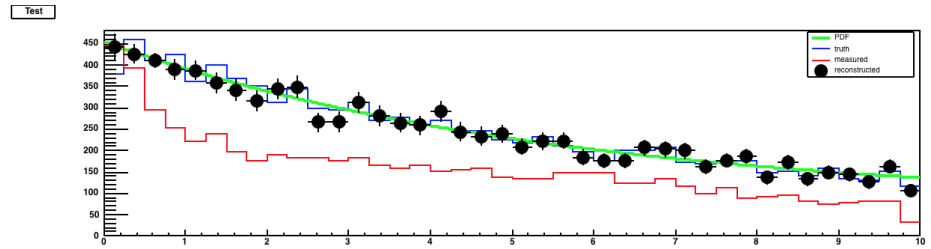
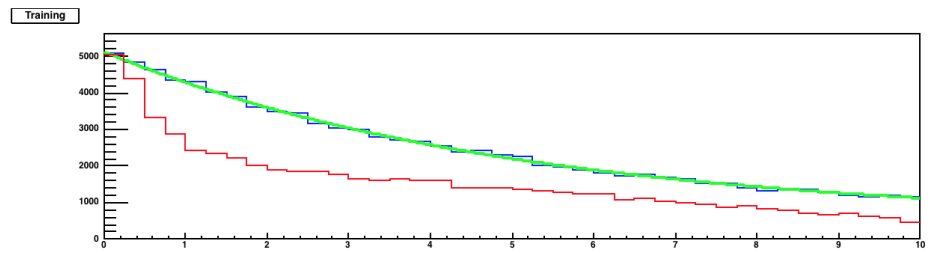
- Rearrange $\vec{v} = R\vec{\mu} + \vec{\beta}$ to get $\vec{\mu} = R^{-1}(\vec{v} - \vec{\beta})$
- Problem: we don't have \vec{v} , we have \vec{n} , the measured data, which is subject to statistical fluxuations.
- We assume n_i is the maximum likelihood estimator for v_i , then solve for the estimator $\hat{\mu} = R^{-1}(\vec{n} - \vec{\beta})$.
- R^{-1} is obtained from R through simple matrix inversion

Iterative Bayesian Method

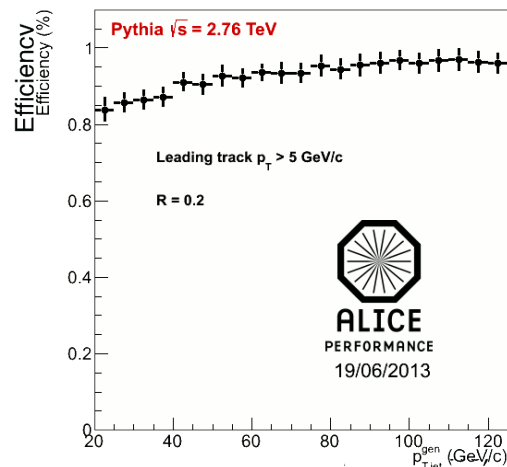
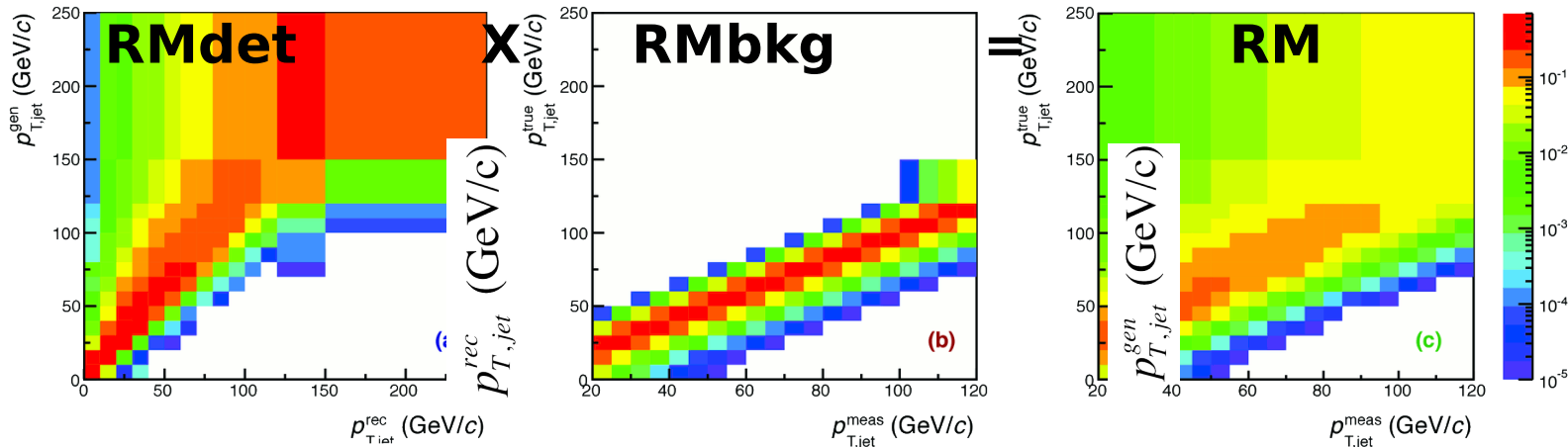
- Using prior knowledge, start with an initial guess for the distribution of true histograms $P^0(\hat{\mu})$
- Use Bayes' Theorem to invert the response matrix $P(\hat{\mu}_i | v_j^{sig}) = \frac{P(v_j^{sig} | \hat{\mu}_i) P^0(\hat{\mu}_i)}{\sum_{l=1}^M P(v_j^{sig} | \hat{\mu}_l) P^0(\hat{\mu}_l)}$
- $\hat{\mu}_i = \frac{1}{\epsilon_i} \sum_{j=1}^N v_j^{sig} P(\hat{\mu}_i | v_j^{sig})$ where ϵ_i is the detector efficiency
- Plug in the newly obtained $P(\hat{\mu}_i | v_j^{sig})$ and $\hat{\mu}_i$ as new priors, then repeat
- Terminate before the wildly oscillating true inverse is reached (usually ~ 4 iterations) to preserve some smoothness

RooUnfold-Bayes

- RooUnfoldTest.cxx
- method = Bayes
- Exponential training and testing



Jets in ALICE: Response Matrix Construction



Anti- k_T R=0.2

$p_{T,track} > 0.15$ GeV/c

$E_{T,cluster} > 0.30$ GeV

$p_{T,track}^{leading} > 5$ GeV/c

(a) **RM_{det}** Detector response matrix

(b) **RM_{bkg}** Background fluctuation matrix

(c) **RM_{tot}** = **RM_{bkg}** × **RM_{det}**

Pb-Pb $\sqrt{s_{NN}} = 2.76$ TeV

0-10% Centrality

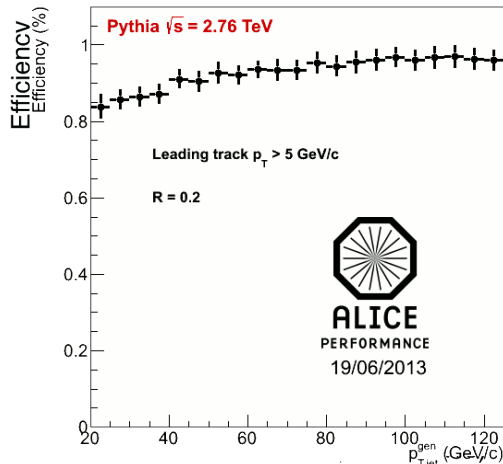
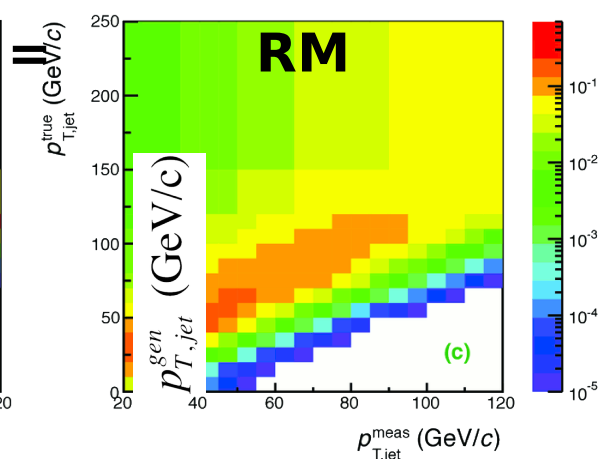
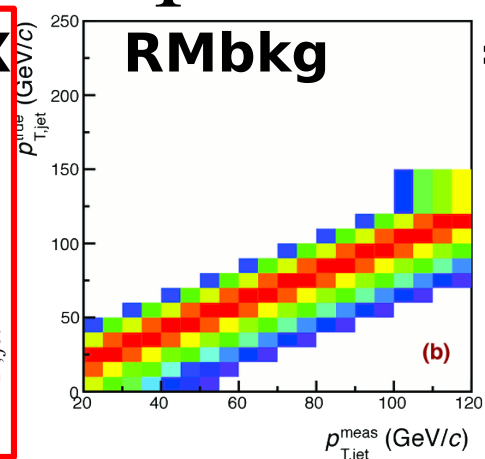
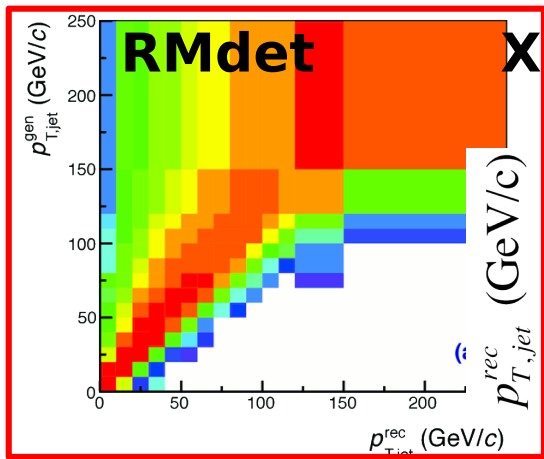


ALICE
PERFORMANCE
15/10/2012

RM_{bkg} and RM_{det} are approximately factorizable

Jets in ALICE: Response Matrix Construction

DETECTOR EFFECT



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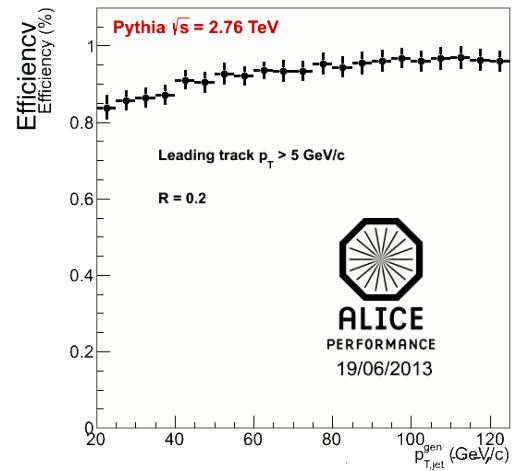
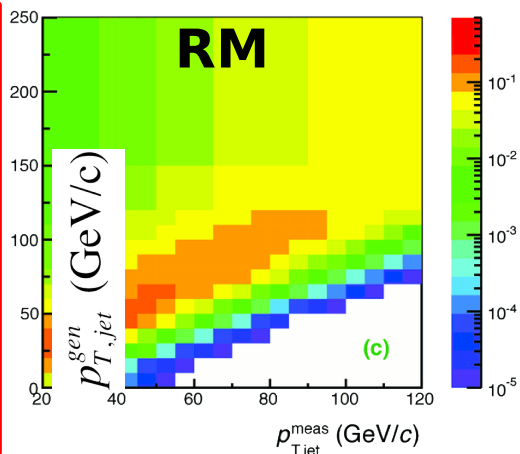
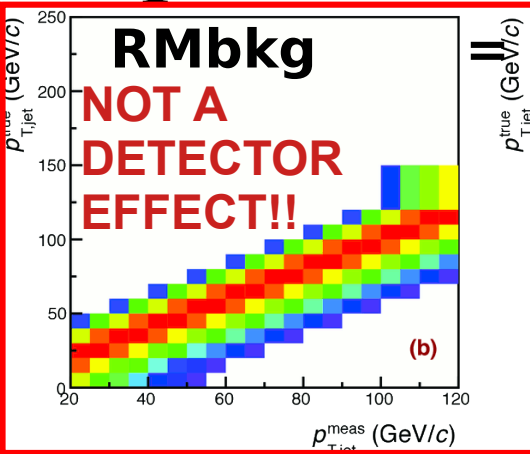
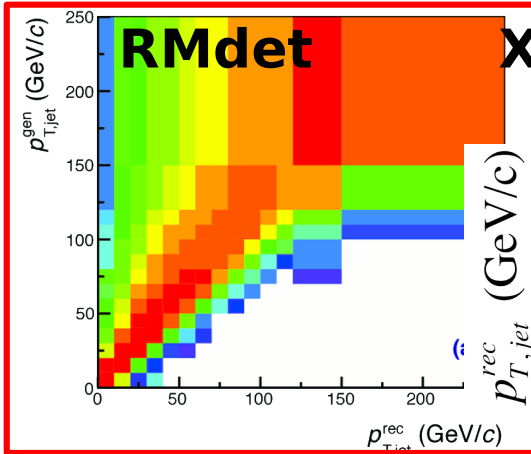


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ALICE
PERFORMANCE
15/10/2012

RM_{bkg} and RM_{det} are approximately factorizable

ALI-PERF-44520

About unfolding...

- d'Agostini (author of Bayesian unfolding algorithm) says you should avoid it if you can
- Necessary when experimental resolution is poor
 - Ex: Single particle spectra $\frac{\sigma_p}{p} \ll w_{bin} \rightarrow$ unfolding unnecessary
 - Ex: Jet spectra $\frac{\sigma_p}{p} \approx w_{bin} \rightarrow$ unfolding necessary
- Algorithm assumes response matrix is correct
 - Matching reconstructed and simulated jets is non-trivial!
- Corrects for multiple experimental effects simultaneously
 - Difficult to disentangle different effects
 - Leads to non-trivial uncertainty correlations between data points due to algorithm
 - May not handle systematic correlations between effects correctly

Mini-summary

- Jet energy resolution is fundamentally large
- Unfolding is complicated, often unstable, and hard
- Construction of response matrix includes several assumptions

Jets in A+A collisions: How to compare to models

Snowmass Accord: Apply the same algorithm to data and your model. Then the measurement and the calculation are the same.

Rivet: Apply the same algorithm to data and your model. Then the measurement and the calculation are the same.

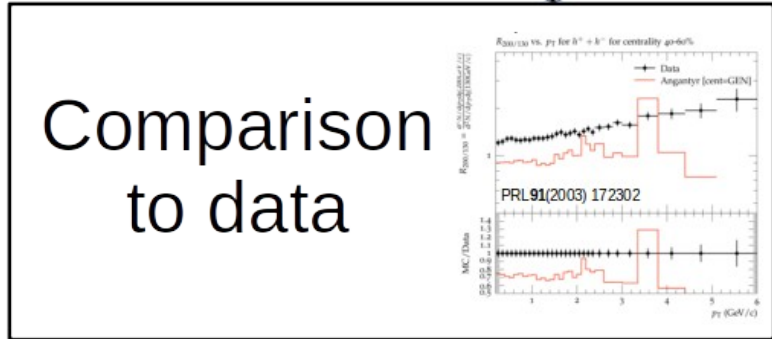
What is Rivet?



HepMC

HEPData

Rivet

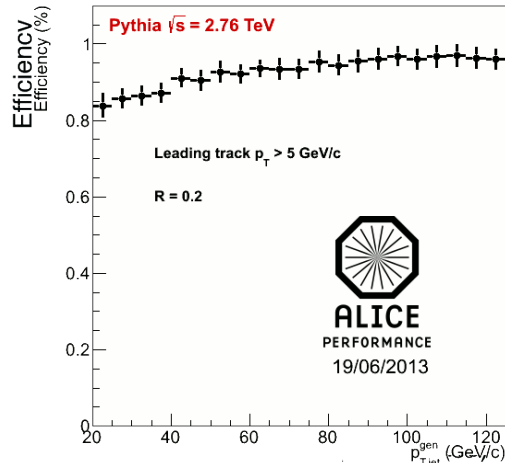
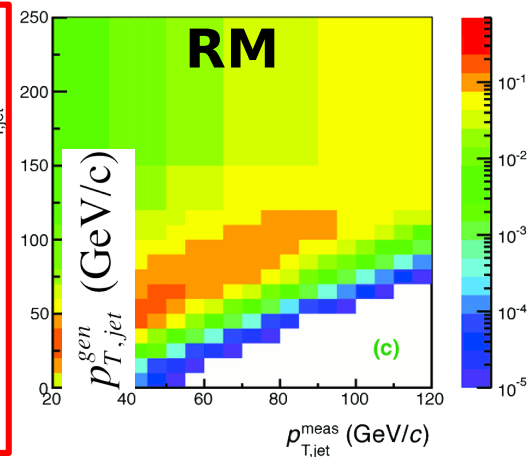
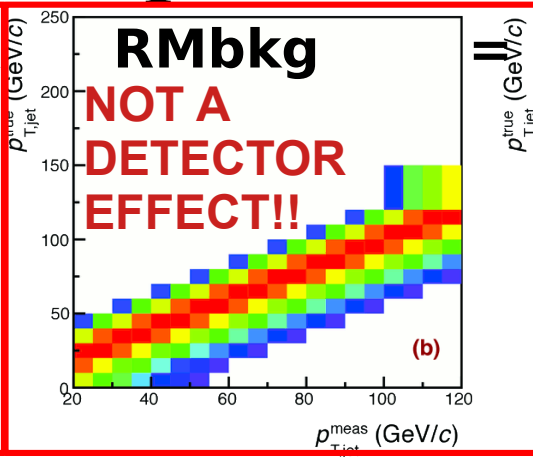
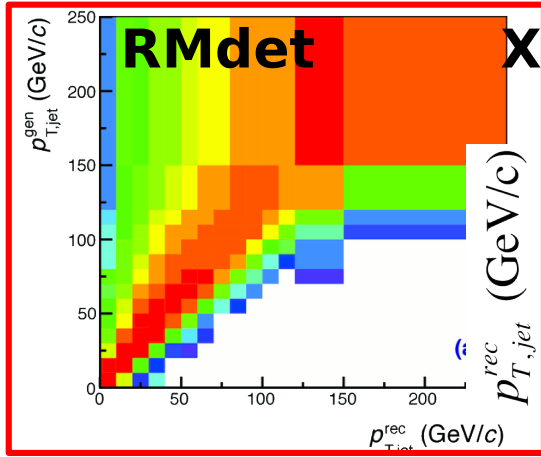


Why use Rivet?

- Facilitates comparisons between Monte Carlos and data
- It's not that hard
- It preserves analysis details

Jets in ALICE: Response Matrix Construction

DETECTOR EFFECT



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ALICE
PERFORMANCE
15/10/2012

RM_{bkg} and RM_{det} are approximately factorizable

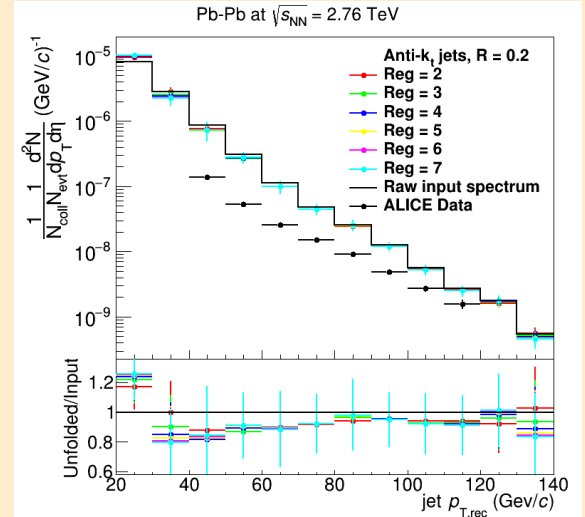
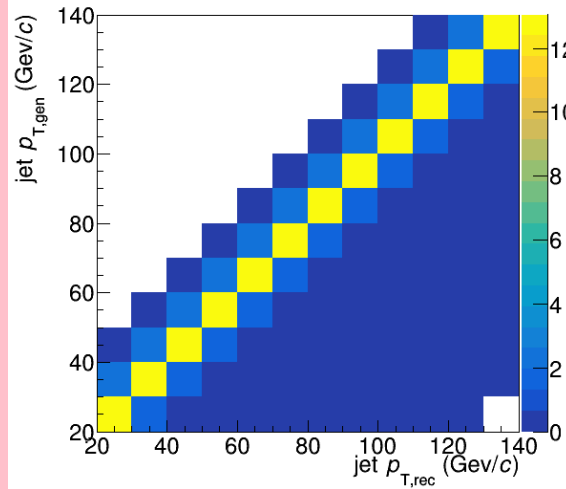
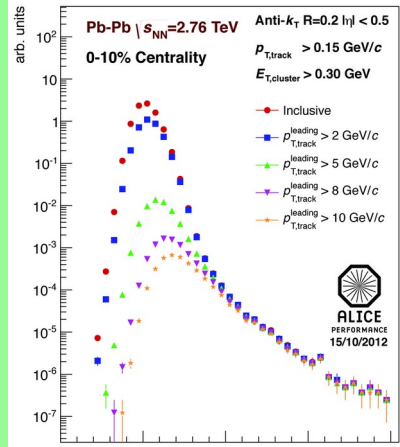
Analysis steps: Full Monte Carlo

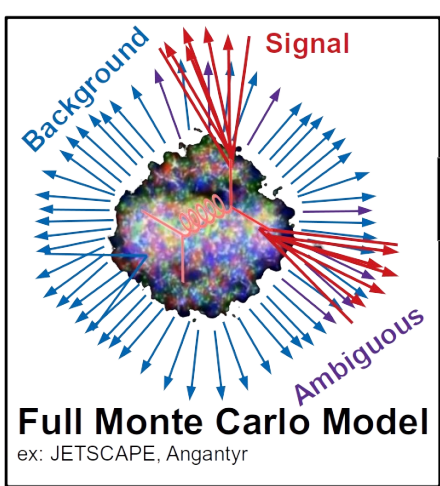
Particles

Jet finding algorithm

Jet candidates

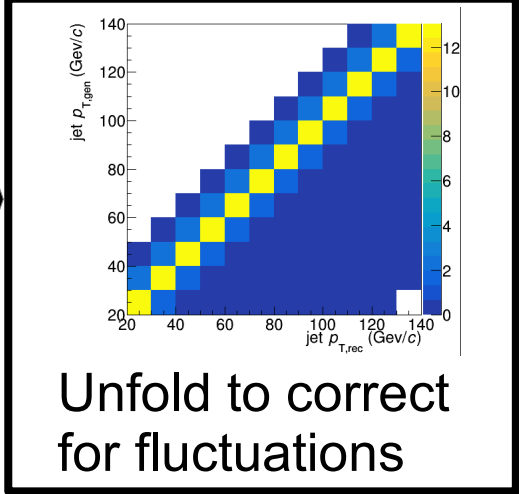
Background subtraction





HepMC

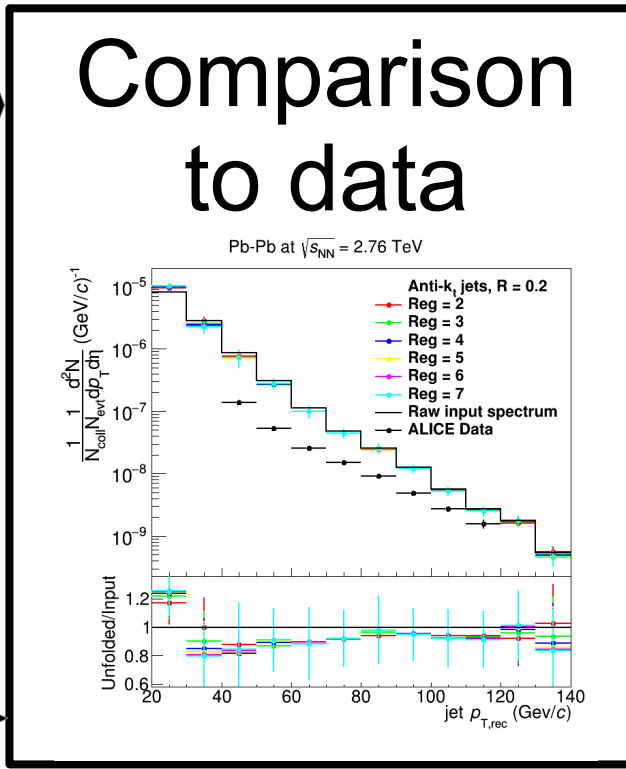
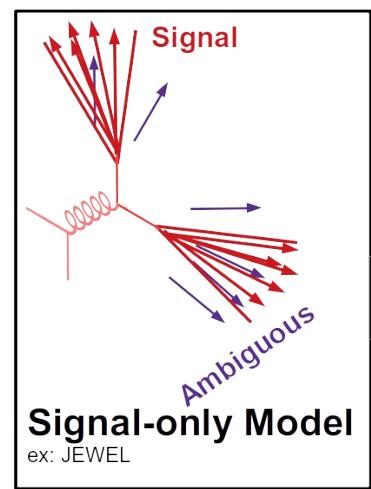
Rivet



HEPData

HepMC

Rivet



Mini-summary

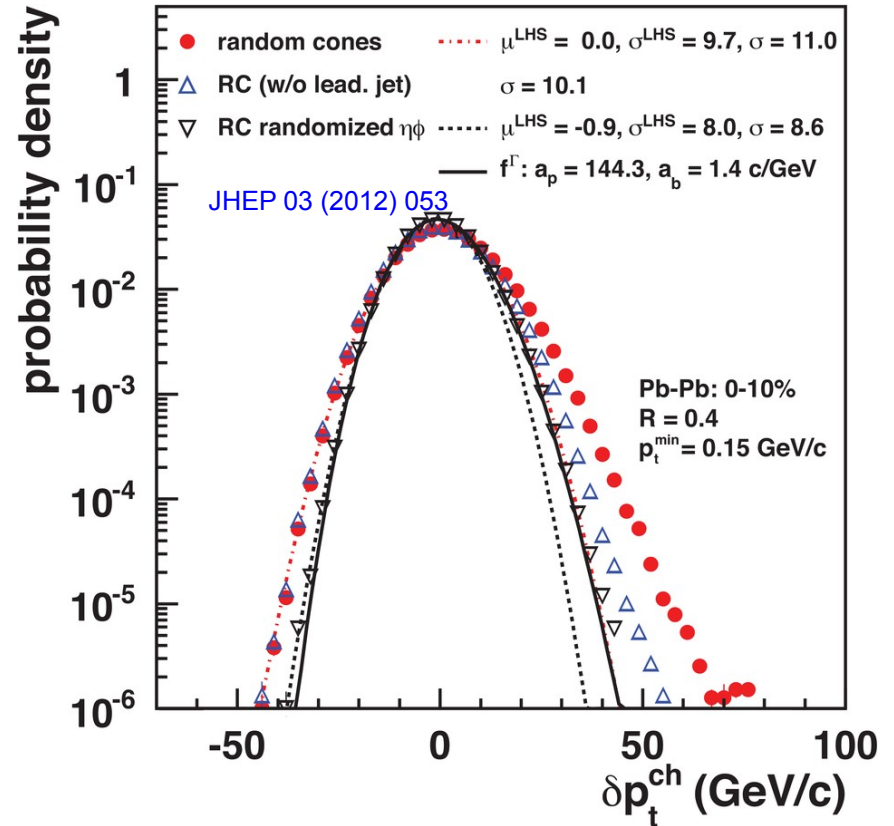
- Experimental techniques can bias measurement in subtle ways
 - Background subtraction
 - Kinematic cuts
 - Choice of jet finder, R
 - Centrality determination
 - Technique for finding reaction plane
- Unclear how these influence the measurement
- Safest to do the same analysis on data and model
 - But unfolding is necessary in a full Monte Carlo model!

Backup

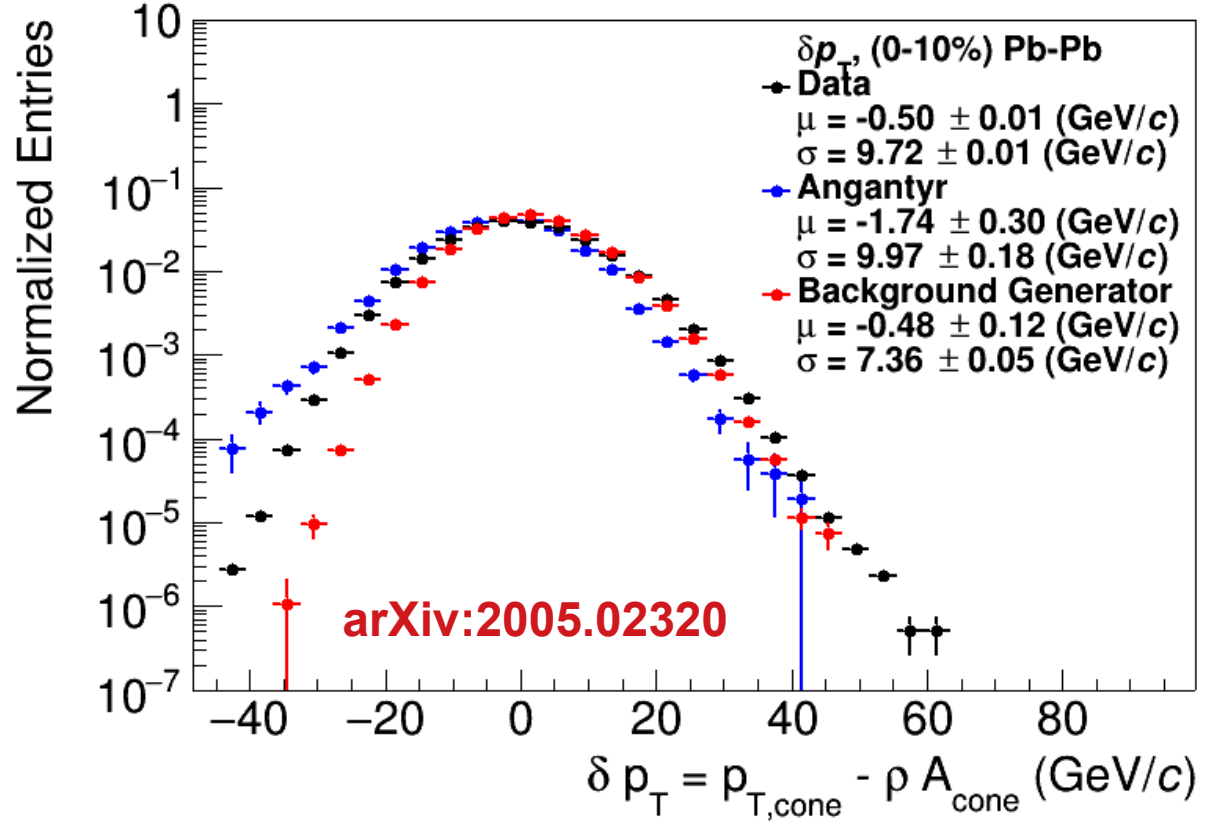
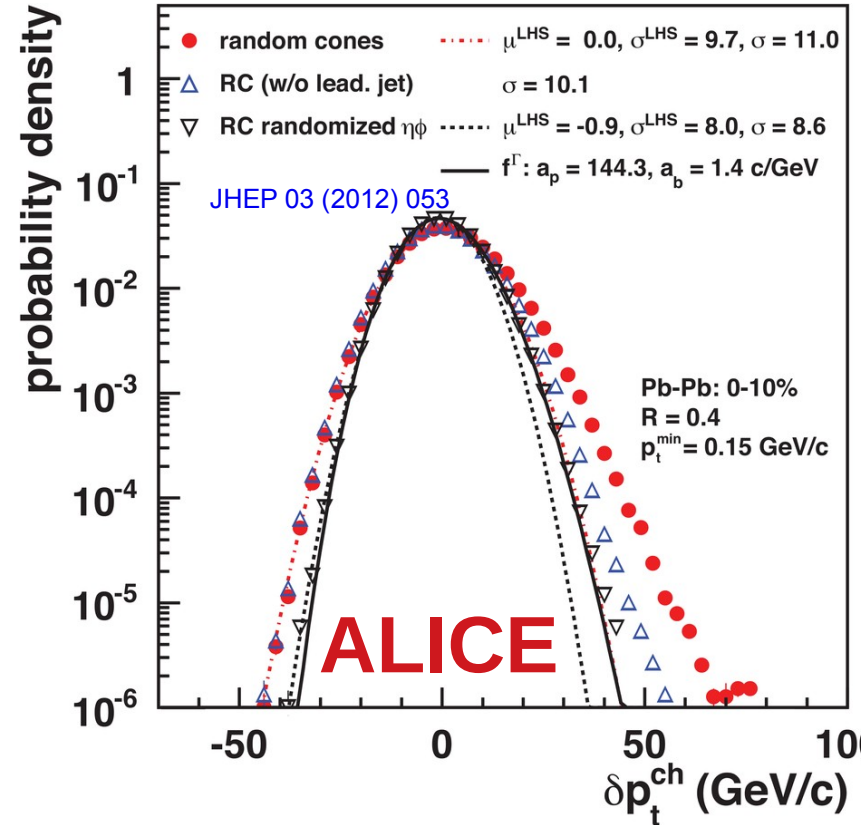
Random cones in ALICE

- Estimate ρ
 - k_T jet finder \rightarrow jet candidates
 - $\rho = \text{Median}(p_T/A)$
- Draw Random cone

$$\delta p_T = p_T^{\text{reco}} - \rho A$$



Random cones



Shape of width of the distribution

Single particle spectra

$$f_{\Gamma}(p_T, p, b) = \frac{b}{\Gamma(p)} (b p_T)^{p-1} e^{-bx}$$

$$\frac{dN}{dy} \propto f_{\Gamma}(p_T, 2, b) = b^2 p_T e^{-k p_T}$$

$$\mu_{p_T} = \frac{p}{b}, \sigma_{p_T} = \frac{\sqrt{p}}{b}$$

Tannenbaum, PLB(498),1-2,Pg.29-34(2001)

Σp_T of N particles \rightarrow N-fold convolution:

$$f_N(p_T, p, b) = f_{\Gamma}(p_T, Np, b) \quad \frac{dp_T^{total}}{dy} \propto f_N(p_T, Np, b)$$

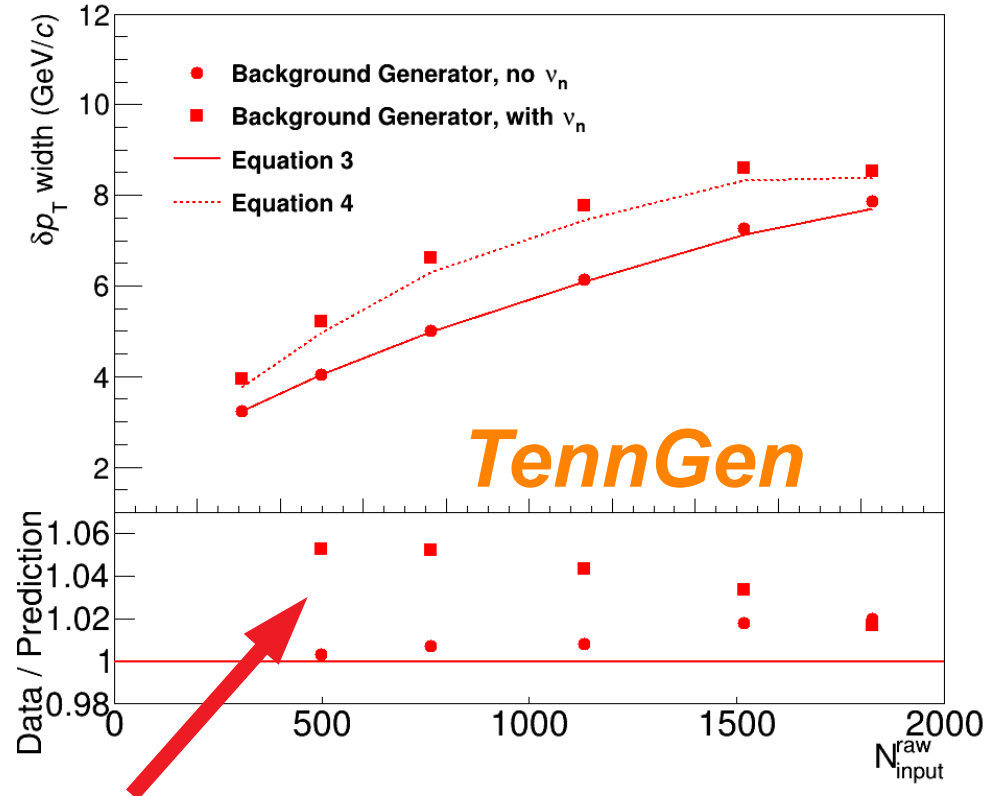
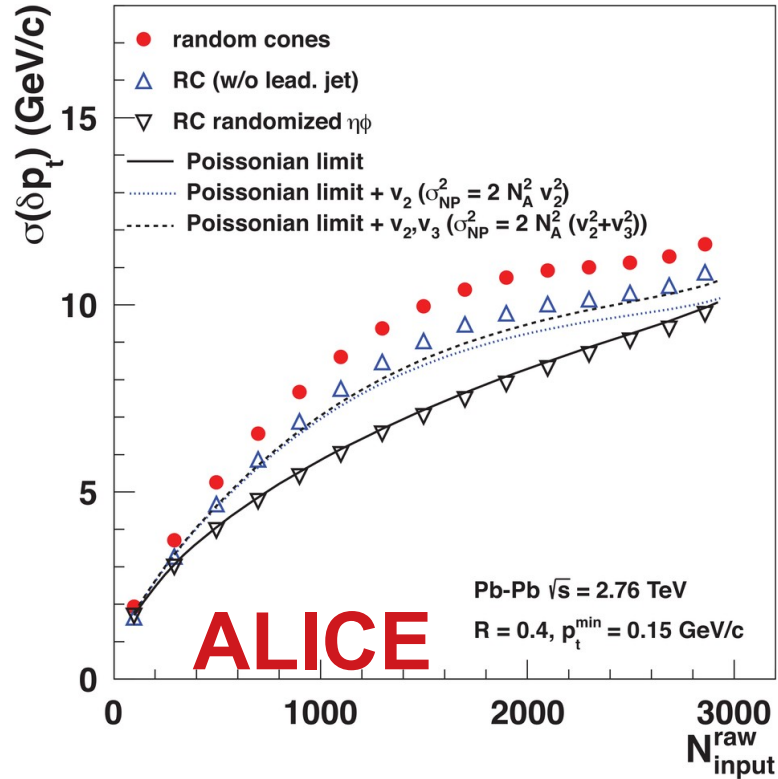
$$N = \frac{N_{total}}{A_{total}} \pi R^2 \quad \mu_{total} = \frac{Np}{b} = N \mu_{p_T}, \sigma_{total} = \frac{\sqrt{Np}}{b} = \sqrt{N} \sigma_{p_T}$$

Add Poissonian fluctuations in N: $\sigma_{total} = \sqrt{N \sigma_{p_T}^2 + N \mu_{p_T}^2}$

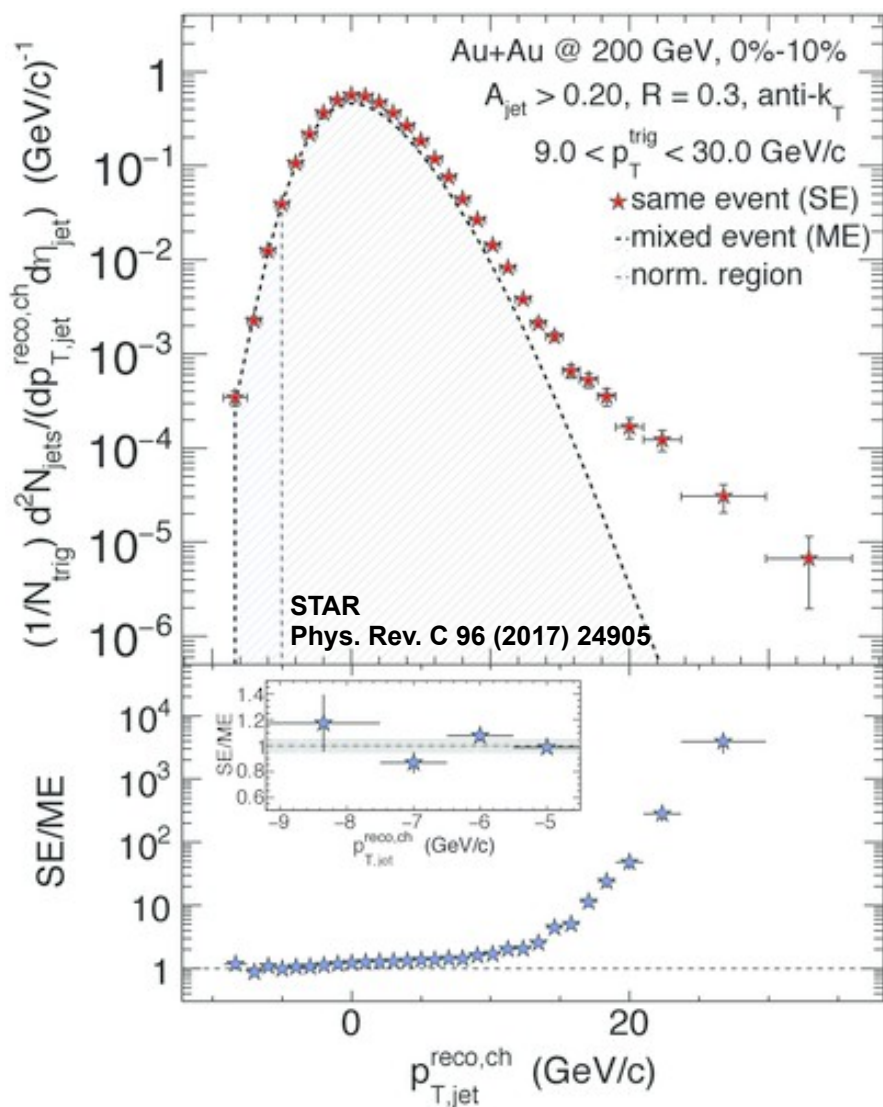
Add non-Poissonian fluctuations in N due to flow

$$\sigma_{total} = \sqrt{N \sigma_{p_T}^2 + (N + 2 \sum_n v_n^2) \mu_{p_T}^2}$$

Width vs multiplicity



Small deviations



Mixed events

- Gets background up to a normalization factor
- Good agreement with the data... but 20% discrepancies still within uncertainties
- In measurement with background suppressed (h-jet correlations)
- Did not see such agreement at the LHC

Shape of width of the distribution

Single particle spectra

$$f_{\Gamma}(p_T, p, b) = \frac{b}{\Gamma(p)} (b p_T)^{p-1} e^{-b p_T}$$

$$\frac{dN}{dy} \propto f_{\Gamma}(p_T, 2, b) = b^2 p_T e^{-k p_T}$$

$$\mu_{p_T} = \frac{p}{b}, \sigma_{p_T} = \frac{\sqrt{p}}{b}$$

Tannenbaum, PLB(498),1-2,Pg.29-34(2001)

Assumes shape

Σp_T of N particles \rightarrow N-fold convolution:

$$f_N(p_T, p, b) = f_{\Gamma}(p_T, Np, b) \quad \frac{dp_T^{total}}{dy} \propto f_N(p_T, Np, b)$$

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$$\text{Add Poissonian fluctuations in N: } \sigma_{total} = \sqrt{N \sigma_{p_T}^2 + N \mu_{p_T}^2}$$

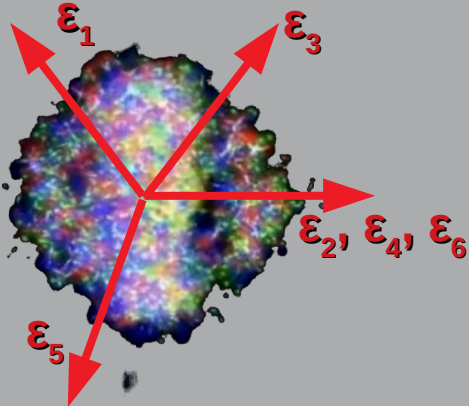
Add non-Poissonian fluctuations in N due to flow

$$\sigma_{total} = \sqrt{N \sigma_{p_T}^2 + (N + 2 \sum_n v_n^2) \mu_{p_T}^2}$$

Assumes uncorrelated number fluctuations

TennGen background generator

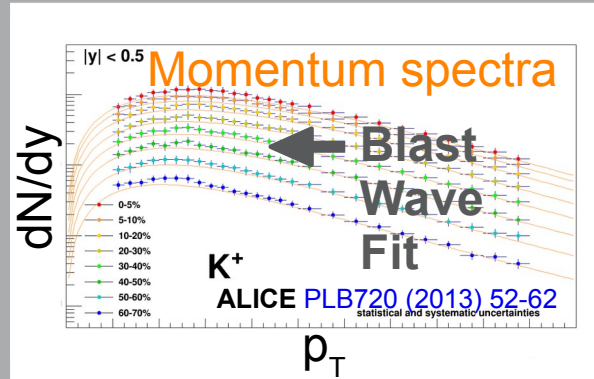
Event properties



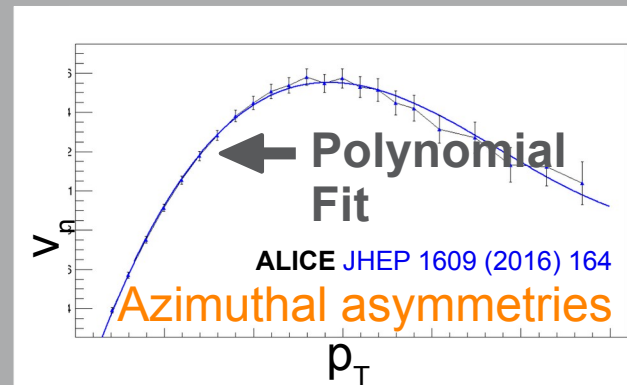
- Even event planes fixed at $\Psi=0$
- Odd planes at random ϕ
- Multiplies from ALICE PRC88 (2013) 044910

**No jets! No resonances
Emulates hydro correlations**

Track properties



→ Random p_T



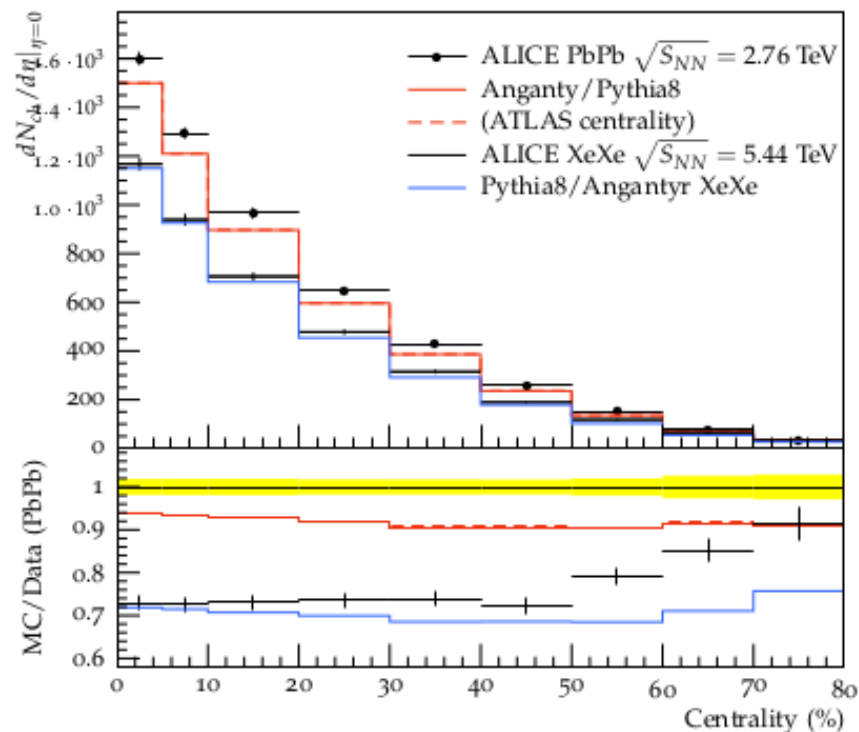
→ v_n
→ Random ϕ

PYTHIA Angantyr

JHEP (2018) 2018: 134

- Based on PYTHIA 8
Sjöstrand, Mrenna & Skands,
JHEP05 (2006) 026
Comput. Phys. Comm. 178 (2008) 852.
- Based on Fritiof & wounded nucleons
- N-N collisions w/fluctuating radii
→ fluctuating σ

**Lots of jets! And resonances!
No hydrodynamics, no jet quenching**



Area-based background subtraction

Cacciari & Salam, [PLB659:119–126,2008](#)

Particles, clusters

k_T algorithm

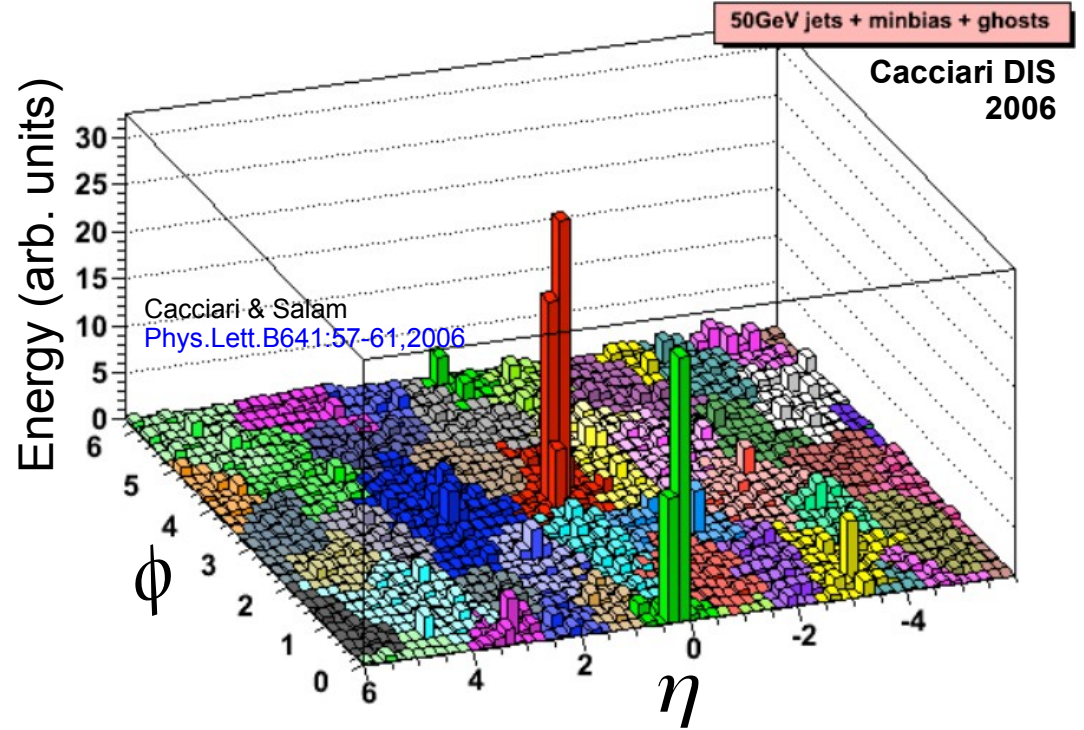
$$k_T = p_T, \Delta R_{ij} = \sqrt{(\eta_i - \eta_j)^2 + (\phi_i - \phi_j)^2}$$

- For all i, j calculate:
$$d_{ij} = \min(p_{T,i}^2, p_{T,j}^2) \Delta R_{ij}^2$$
 - $$d_{iB} = p_{T,i}^2$$
 - Combine smallest d_{ij} .
If d_{iB} smallest, $d_{iB} \rightarrow$ jet
- Repeat until no particles left

Jet candidates

Median $\rho = p_T / A$

$$p_T^{jet} = p_T^{reco} - \rho_{median} A^{jet}$$



Theoretical calculations

Factorization theorem

- Assumption: Parton distribution functions, perturbative cross section, fragmentation function factorize
- What people really mean by “perturbatively calculable”
 - D and f are *explicitly non-perturbative!*
 - D is for *parton $c \rightarrow$ hadron h*
Not what is experimentally measured
- Most theories for jet quenching modify fragmentation function D

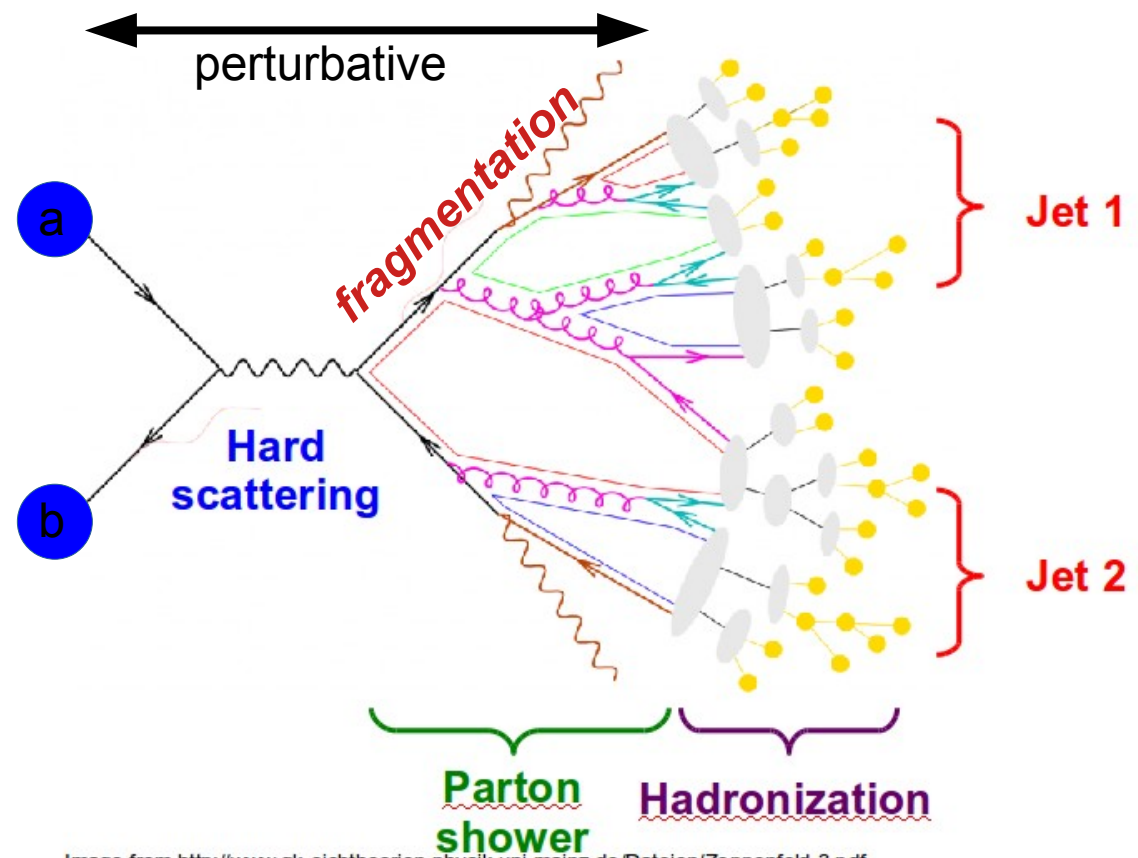


Image from <http://www.gk-eichtheorien.physik.uni-mainz.de/Dateien/Zeppenfeld-3.pdf>

$$\frac{d^3 \sigma^h}{dy d^2 p_T} = \frac{1}{\pi} \int d x_a \int d x_b f_a^A(x_a) f_b^B(x_b) \frac{d \sigma_{ab \rightarrow cX}}{d \hat{t}} \frac{D_c^h(z)}{z}$$

Jet finders